

# Experimental Search for Spin-Dependent (and Independent) Forces in the Sub-mm range

**Josh Long**

**Simon Kelly, Evan Weisman, Trevor Leslie, Andrew Peckat**  
*Indiana University, Bloomington*



**Parameterization and Existing Limits**

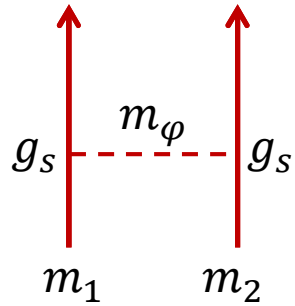
**Experimental Approach**

**Projected Sensitivity**

**Recent Data and Current Sensitivity (spin-independent)**

**Spin-Polarized Test Mass R&D**

# Parameterization



## Yukawa

$$V(r) = -G \frac{m_1 m_2}{r} [1 + \alpha e^{-r/\lambda}]$$

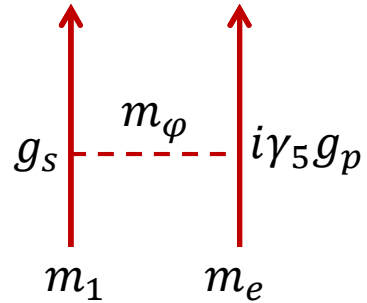
$$\lambda = \hbar/m_\phi c = \text{range}$$

$\alpha$  = strength relative to gravity

$\sim "V_1"$  [1]

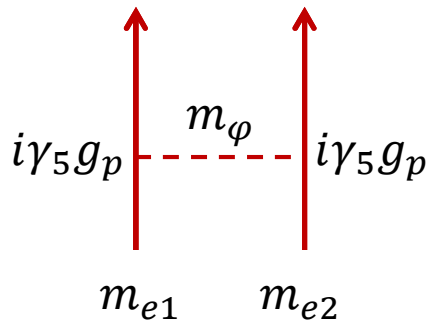
# Spin-Dependent Interactions

## Monopole-dipole [2]



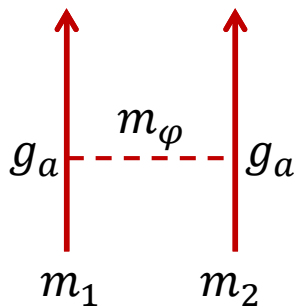
$$V(r) = -g_p g_s \frac{\hbar^2}{8\pi m_e} \hat{\sigma} \cdot \hat{r} \left( \frac{1}{r\lambda} + \frac{1}{r^2} \right) e^{-r/\lambda} \quad \sim "V_9 + V_{10}" \text{ (P, T-odd)}$$

## Dipole-dipole



$$V(r) = \frac{g_p g_p \hbar^3}{4 m_{e1} m_{e2} c} [(\sigma_1 \cdot \sigma_2) \left( \frac{1}{\lambda r^2} + \frac{1}{r^3} \right) - (\sigma_1 \cdot \hat{r})(\sigma_2 \cdot \hat{r}) \left( \frac{1}{\lambda^2 r} + \frac{3}{\lambda r^2} + \frac{3}{r^3} \right)] e^{-r/\lambda} \quad \sim "V_3"$$

## $V_2$



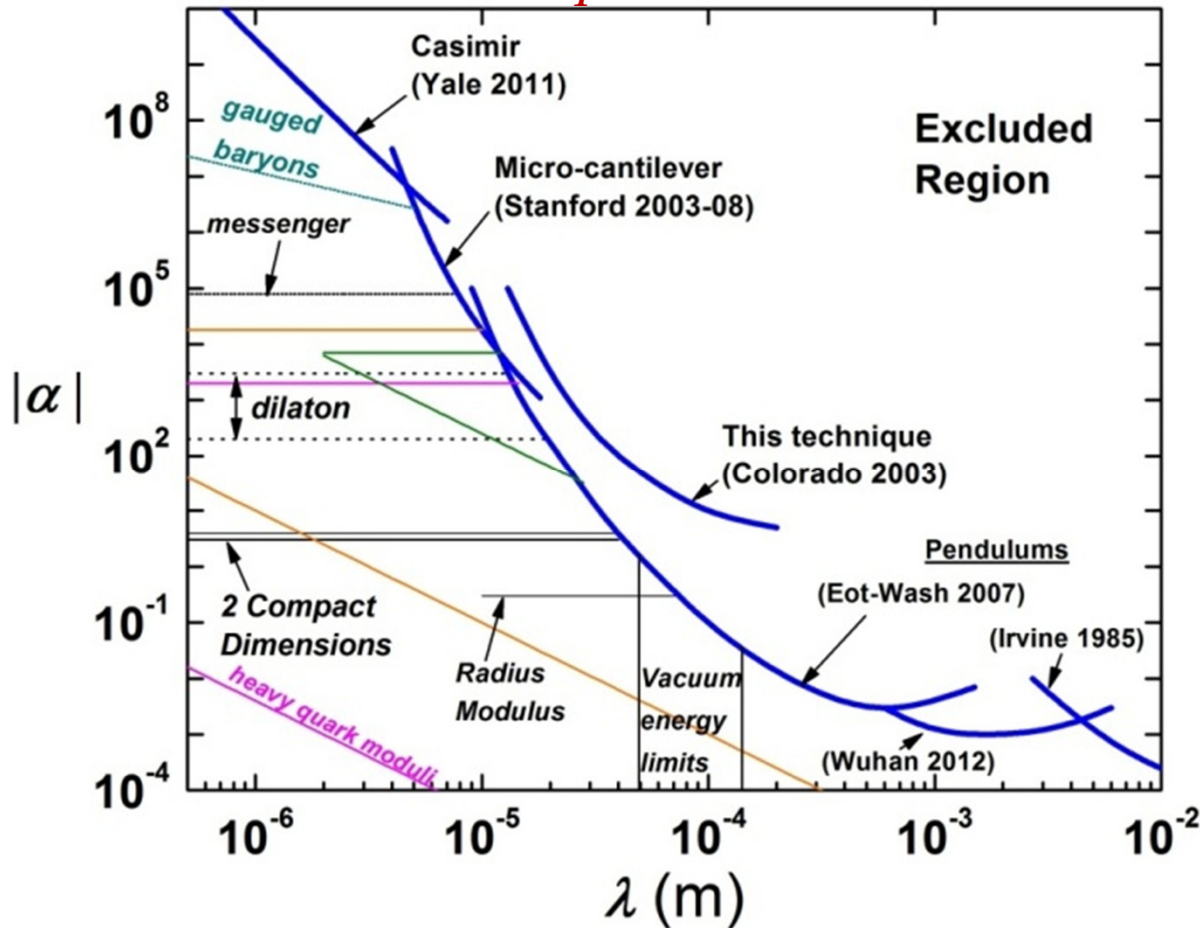
$$V(r) = -g_a g_a \frac{\hbar c}{r} (\sigma_1 \cdot \sigma_2) e^{-r/\lambda}$$

[1] B. Dobrescu and I. Mocioiu, J. High Energy Phys. 0611, 005 (2006)

[2] J. E. Moody and F. Wilczek, Phys. Rev. D 30, 130 (1984)

# Direct Experimental Limits and Predictions

$V_1$  - Yukawa



## Experimental limits:

Irvine, Wuhan, Eot-Wash, Yale = torsion pendulum experiments

Stanford = AFM-type experiment

Limits still allow forces 1 million times stronger than gravity at 5 microns

## Theoretical predictions:

“Large” extra dimensions

Vacuum energy: prediction from new field which also keeps cosmological constant small

Moduli, dilatons: new particles motivated by string models

Irvine: J. Hoskins et al., PRD 32 (1985) 3084

Wuhan: S-Q. Yang et al., PRL 108 (2012) 081101

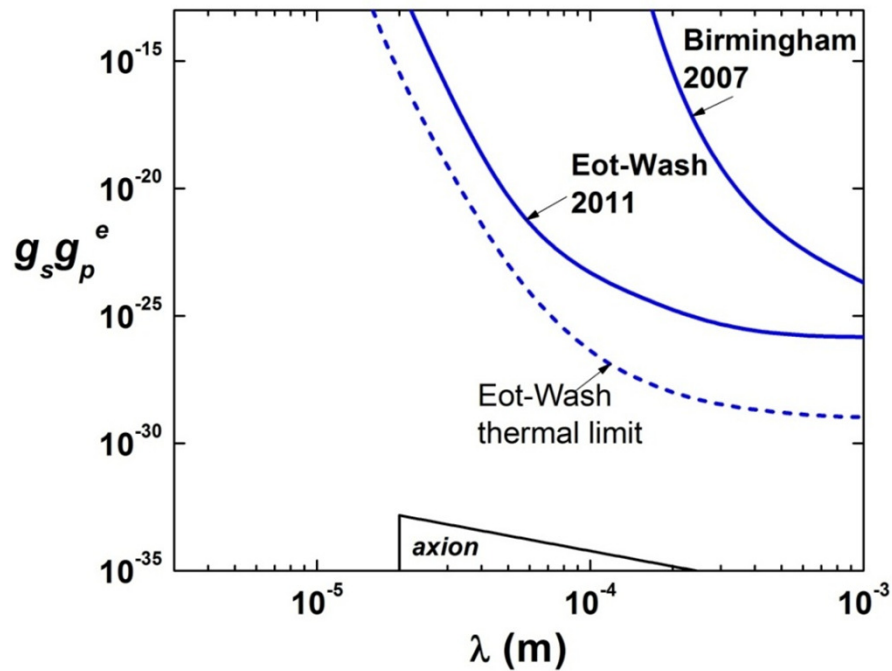
Eot-Wash: D. Kapner et al., PRL 98 (2007) 021101

Stanford: A. Geraci et al., PRD 78 022002 (2008)

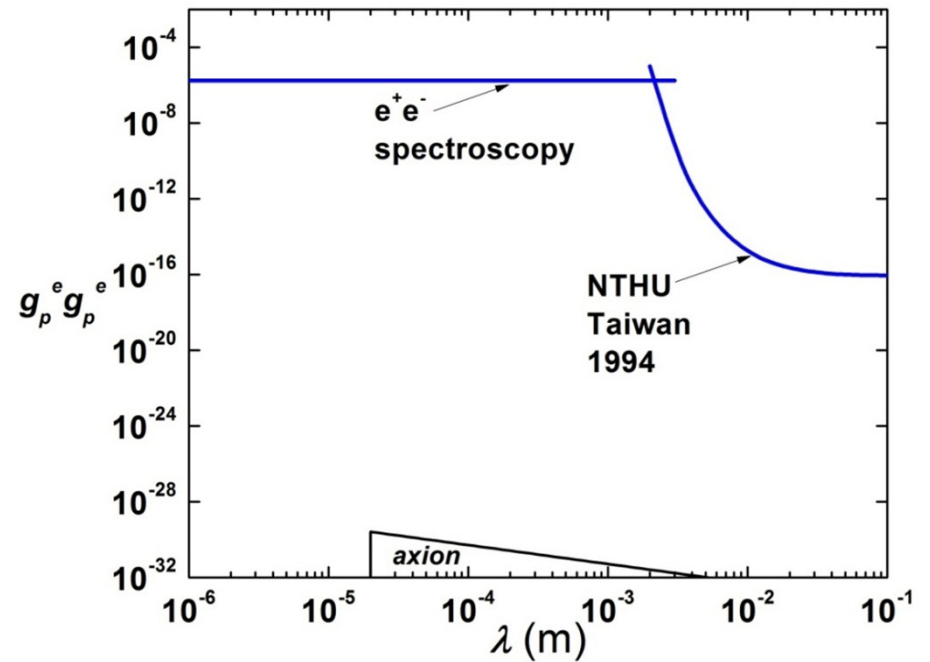
Yale: A. Sushkov et al., PRL 107 (2011) 171101 **[APS: H12 - 1]**

# Direct Experimental Limits and Predictions

## $V_{9+10}$ - monopole-dipole



## $V_3$ - dipole-dipole



Eot-Wash: S. Hoedl et al., PRL 106 (2011) 041801

NTHU: W-T. Ni et al., Physica B 194 (1994) 153

$e^+e^-$ : A. Mills, PRA 27 (1983) 262;

M. Ritter et al., PRA 30 (1984) 1331

# Experimental Approach

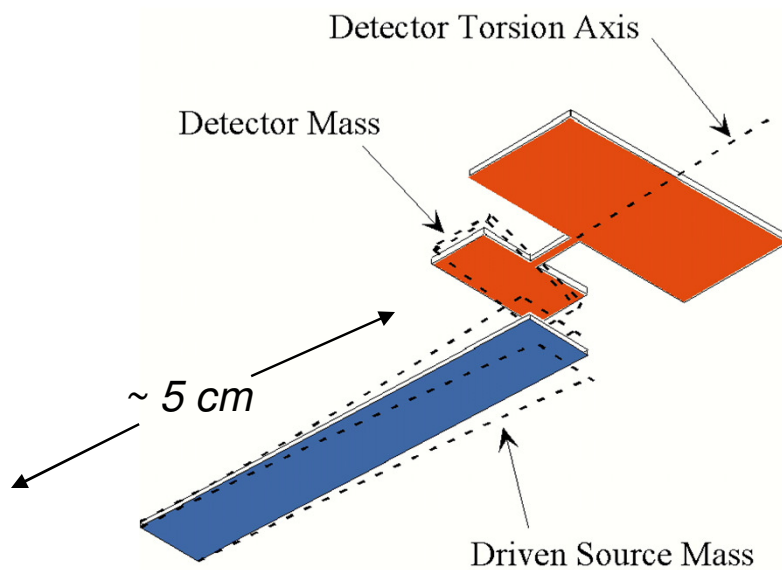
**Planar Geometry - null for  $1/r^2$**

**Resonant detector with source mass driven on resonance**

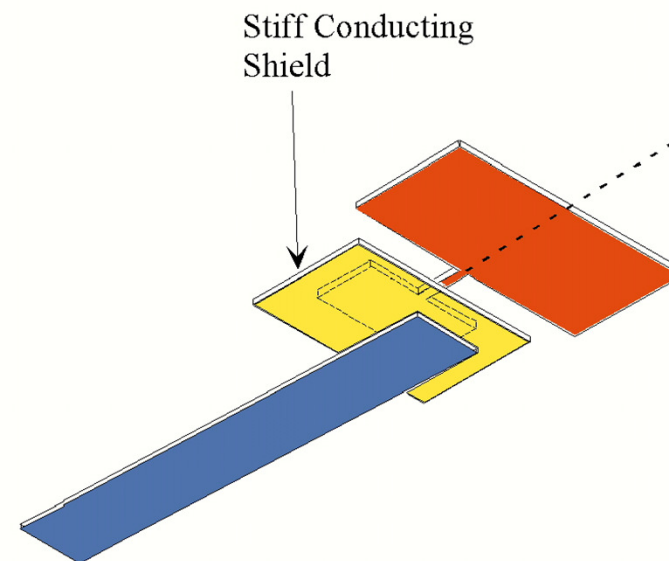
**1 kHz operational frequency - simple, stiff vibration isolation**

**Double-rectangular torsional detector: high Q, low thermal noise**

**Stiff conducting shield for background suppression**



***Source and Detector Oscillators***



***Shield for Background Suppression***

# Central Apparatus

Vibration isolation stacks: Brass disks connected by fine wires; soft springs which attenuate at  $\sim 10^{10}$  at 1 kHz (reason for using 1 kHz)

Readout: capacitive transducer and lock-in amplifier

Vacuum system:  $10^{-7}$  torr

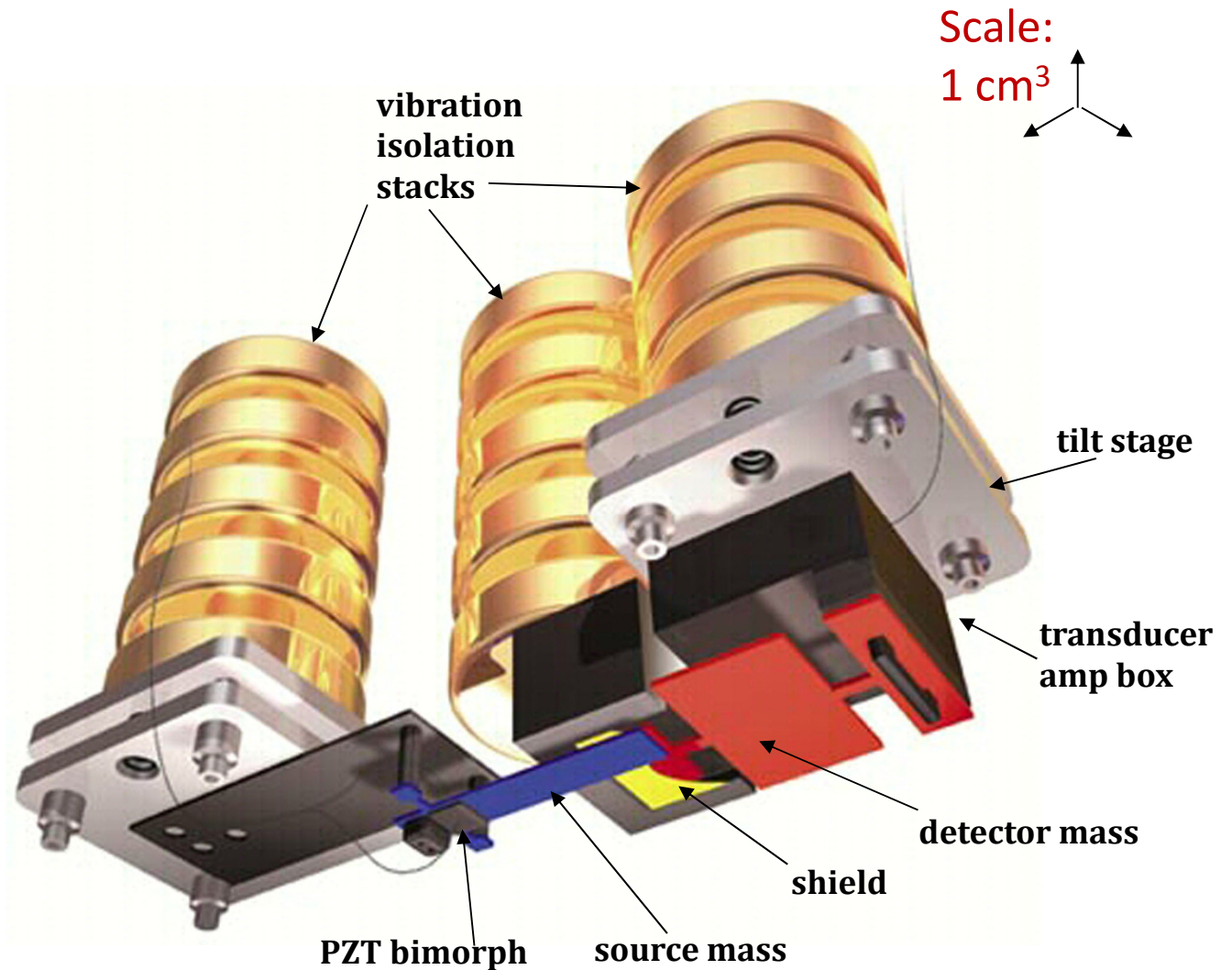
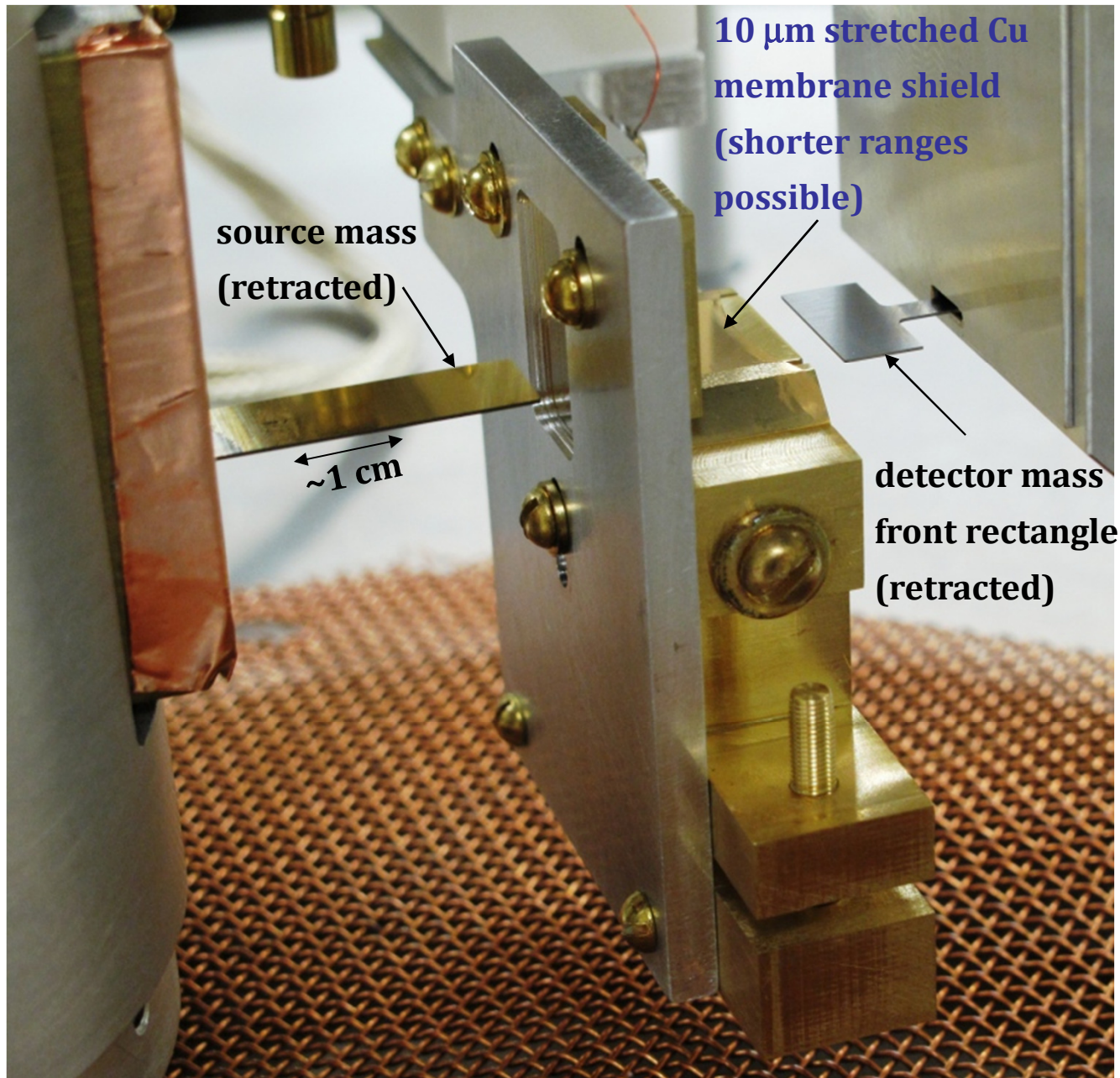


Figure: Bryan Christie ([www.bryanchristie.com](http://www.bryanchristie.com)) for Scientific American (August 2000)

# Interaction Region



## Thinner shield

60  $\mu$ m thick sapphire plate  
replaced by 10  $\mu$ m stretched  
copper membrane

Compliance  $\sim 5$ x better than  
needed to suppress estimated  
electrostatic force

Minimum gap reduced from  
105  $\mu$ m (2003) to 40  $\mu$ m.



## Sensitivity: increase Q and statistics, decrease T

- **Yukawa signal: Force on detector due to Yukawa interaction with source**

$$F_Y(t) \approx 2\pi\alpha G\rho_s\rho_d A_d \lambda^2 \exp(-d(t)/\lambda)[1 - \exp(-t_s/\lambda)][1 - \exp(-t_d/\lambda)] \\ \sim 3 \times 10^{-15} \text{ N (for } \alpha = 1, \rho = 20 \text{ g/cc, } \lambda = 50 \text{ } \mu\text{m)}$$

- **Spin-dependent: Integrate  $V_2, V_3$ , monopole-dipole numerically with:**

-  $\sim \text{cm}^2 \times 100 \text{ } \mu\text{m}$  spin-polarized samples on test masses

-  $n_s = 10^{21}/\text{cc}$  (10% of world record [1])

- **Thermal Noise**

$$F_T = \sqrt{\frac{4kTD}{\tau}} \quad D = \frac{\omega m}{Q}$$

$\sim 3 \times 10^{-15} \text{ N rms}$  (300 K,  $Q = 5 \times 10^4$ , 1 day average)

$\sim 7 \times 10^{-17} \text{ N rms}$  ( 4 K,  $Q = 5 \times 10^5$ , 1 day average)

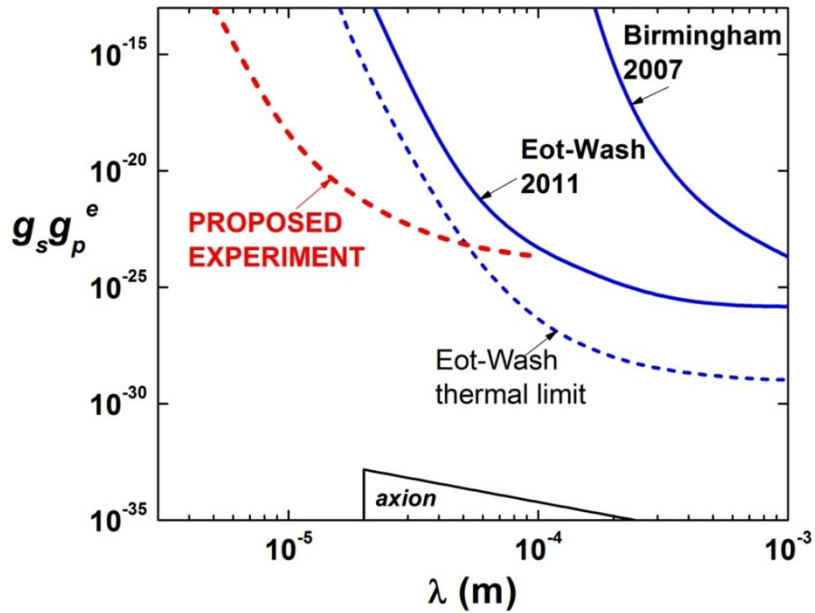
- **Setting SNR = 1 yields**

$$\alpha \sim \frac{1}{\rho^2} \sqrt{\frac{kTm\omega}{Q\tau}} \quad gg \sim \frac{1}{n^2} \sqrt{\frac{kTm\omega}{Q\tau}}$$

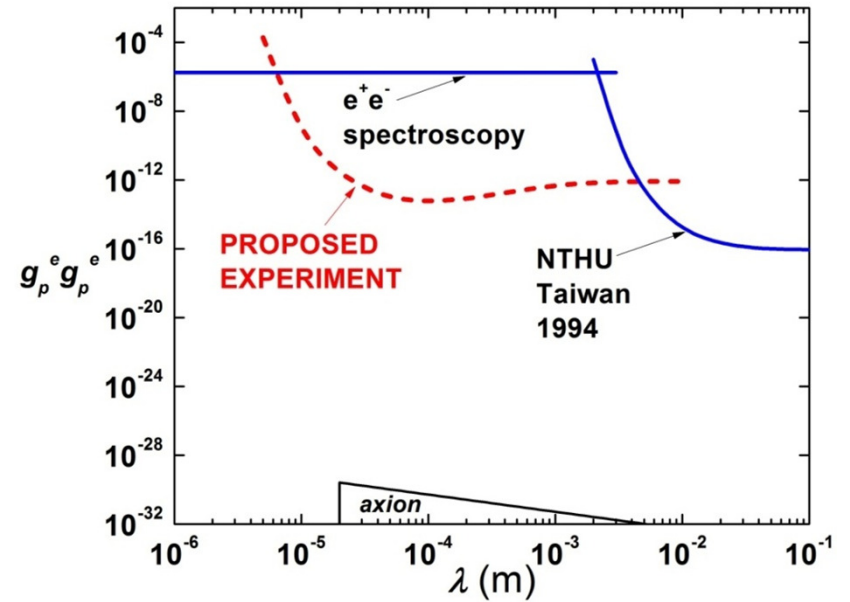
[1] W-T. Ni et al., Physica B 194 (1994) 153

# Projected Sensitivity (Ultimate)

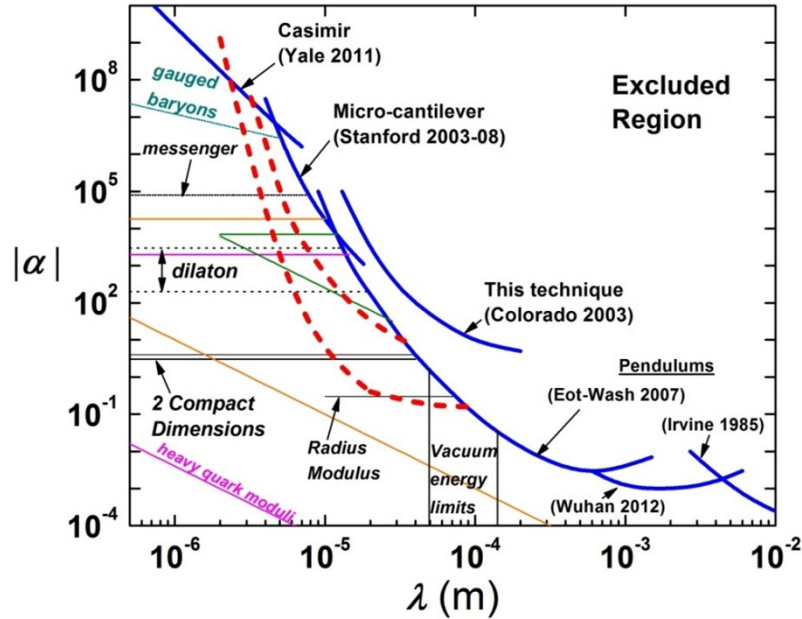
$V_{9+10}$  - monopole-dipole



$V_3$  - dipole-dipole

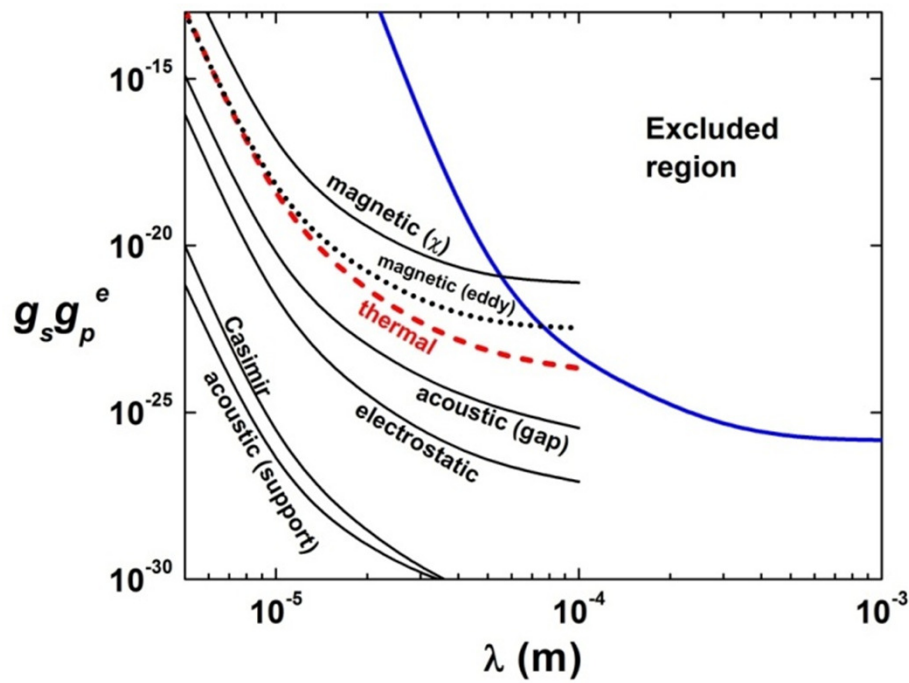


$V_1$  - Yukawa

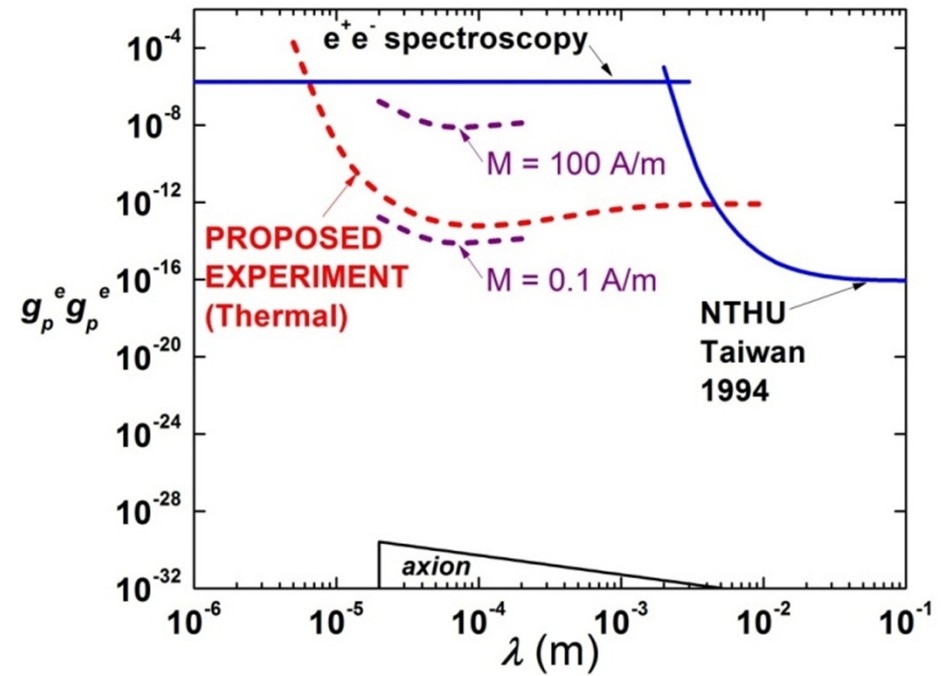


# Other Backgrounds

$V_{9+10}$  - monopole-dipole



$V_3$  - dipole-dipole  
(magnetic contaminants)



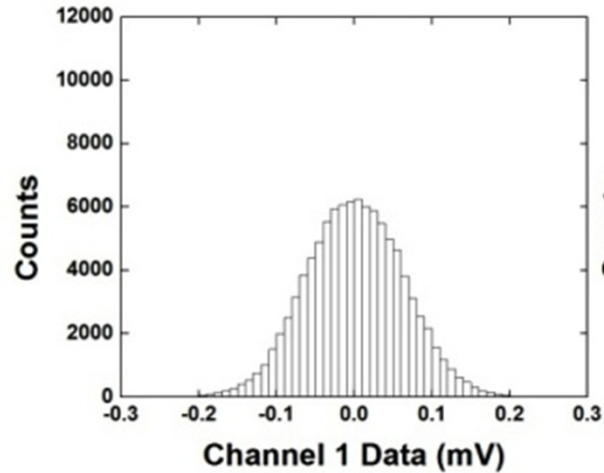
# Spin-Independent Force Measurement Data – March 2012

19 hours on-resonance data collected over 3 days with interleaved diagnostic data

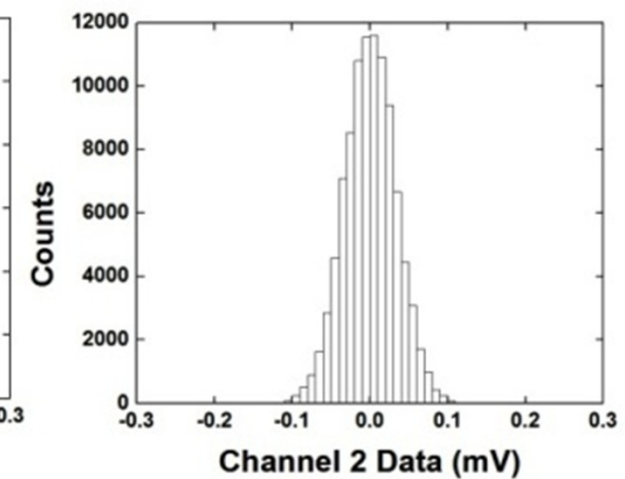
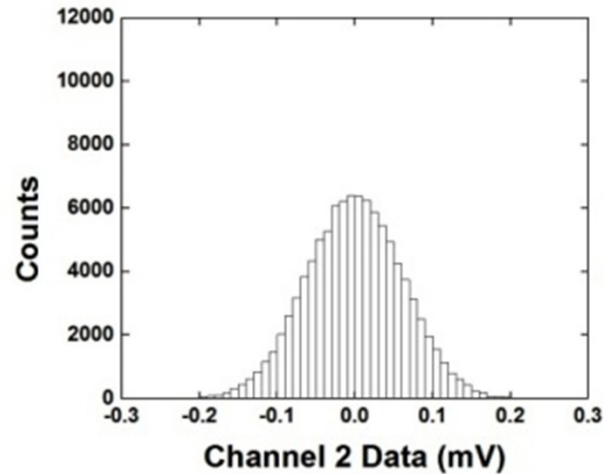
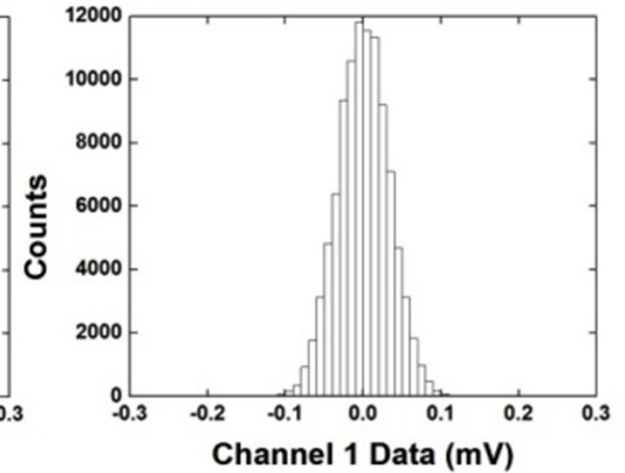
On-resonance: Detector thermal motion and amplifier noise

Off-resonance: amplifier noise

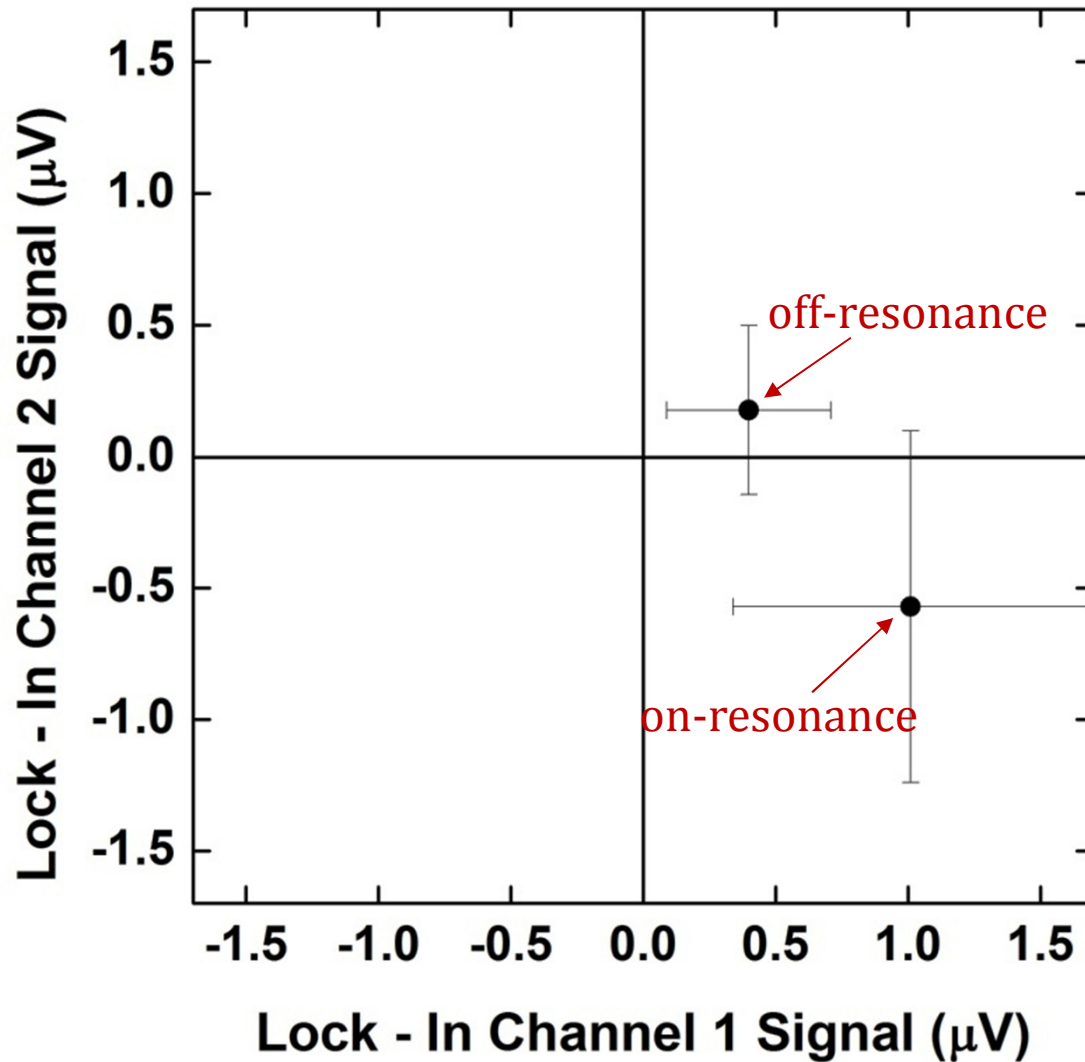
**On Resonance**



**Off Resonance**



## Force Measurement Data - Detail



### Net Signal:

$$V_{\text{on}} - V_{\text{off}} = 0.93 \pm 0.74 \mu\text{V} (1\sigma)$$

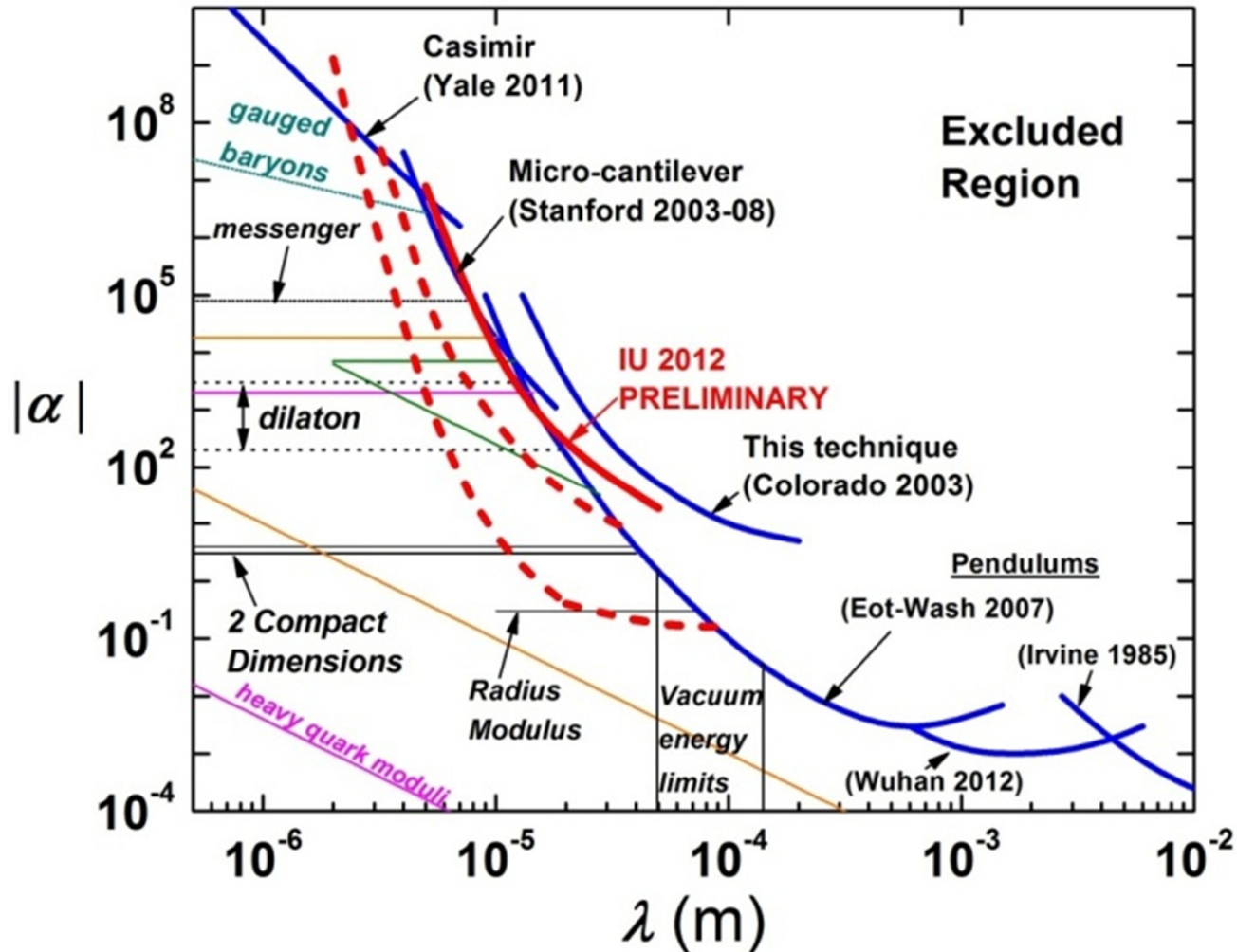
### Force:

$$F = 4.0 \pm 3.2 \text{ fN}$$

### Possible Source:

Detector - *probe* force from  
 $\sim$  nV scale "ground"  
fluctuations on detector  
mass

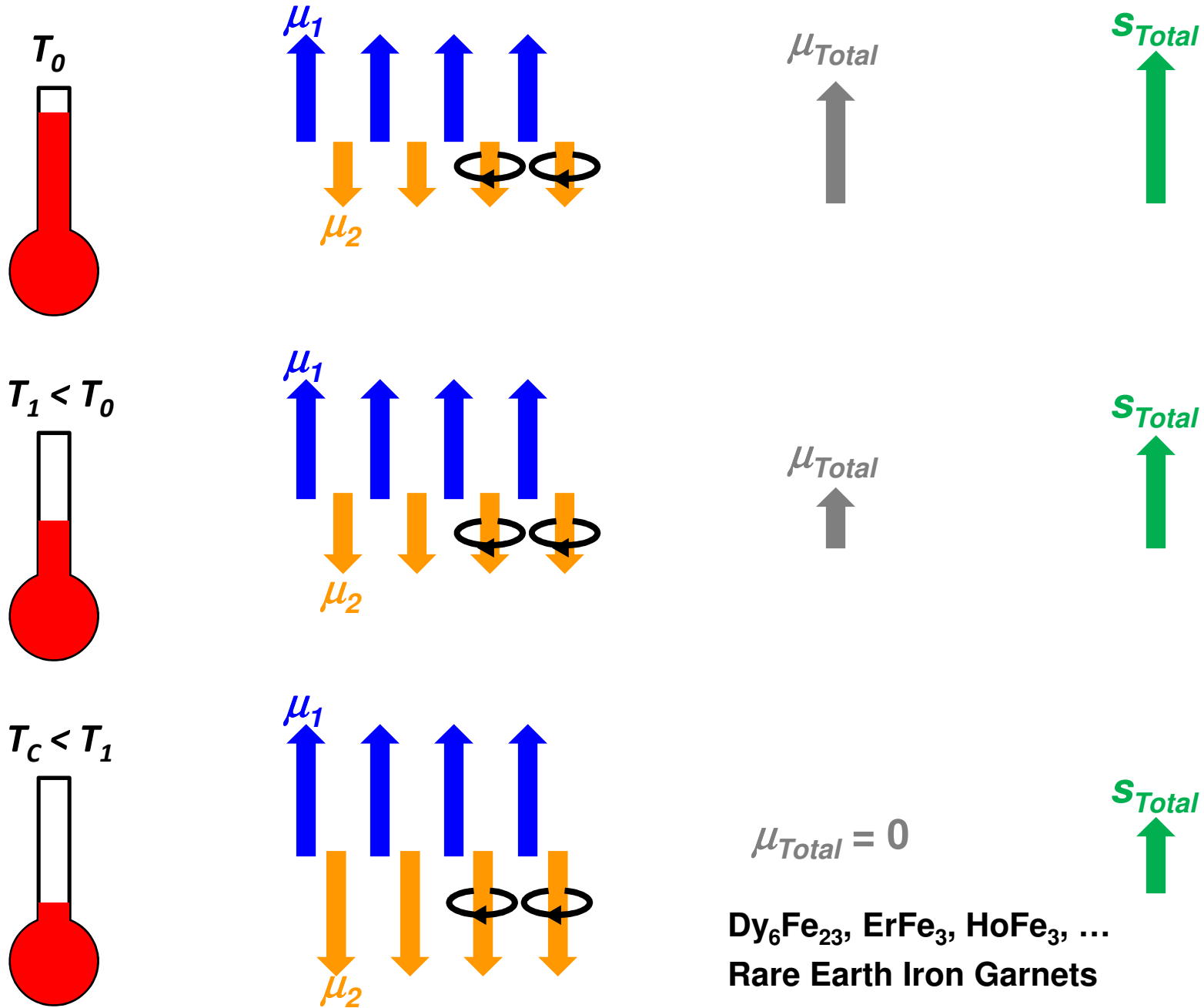
## Current Limits ( $2\sigma$ )



**Minimum gap:  $55 \pm 6 \mu\text{m}$**   
 assumes  $10 \mu\text{m}$  shield flat (optimistic)

**Assumes phase of Yukawa force in direction of maximum signal**  
 (pessimistic; not yet measured)

# Compensated Ferrimagnet

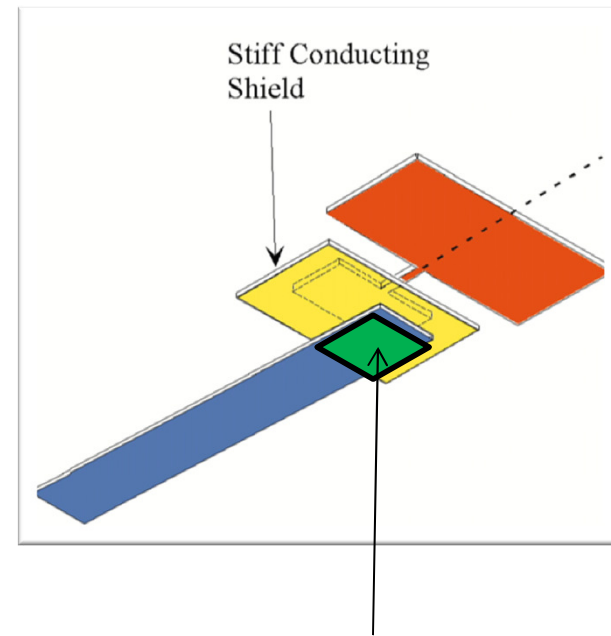
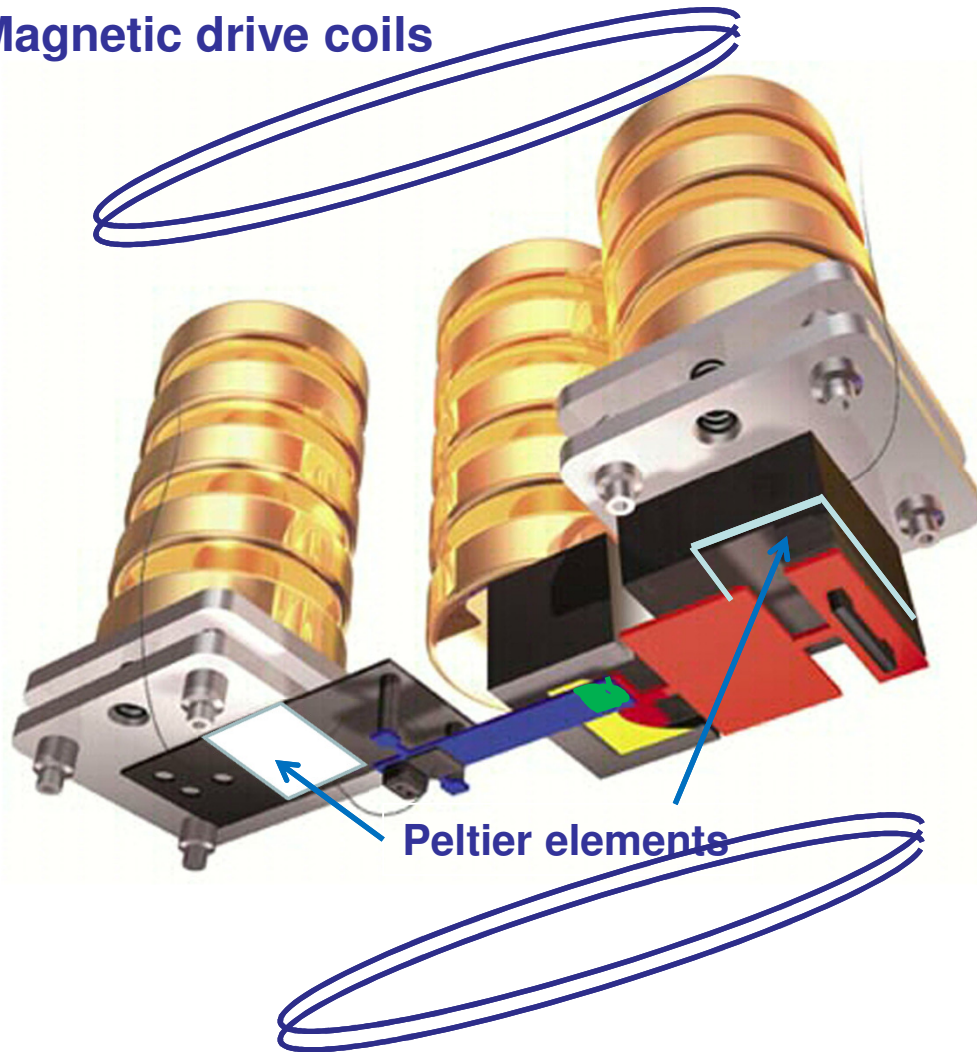


# Spin-Dependent Apparatus

Adapt current apparatus

- Cool to  $T_{\text{comp}}$  with thermoelectric elements
- Drive test masses with AC magnetic field until  $T_{\text{comp}}$  (no magnetization)

Magnetic drive coils

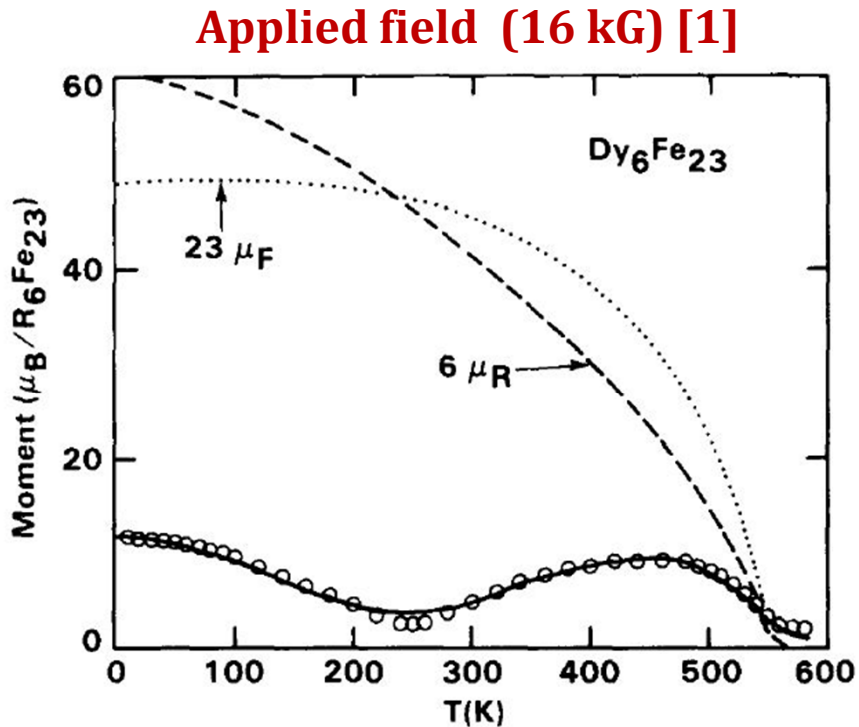


- Dy<sub>6</sub>Fe<sub>23</sub> test mass  
100  $\mu\text{m}$  thick
- Magnetic shield  
100 $\mu\text{m}$ ,  $\mu$ -metal?

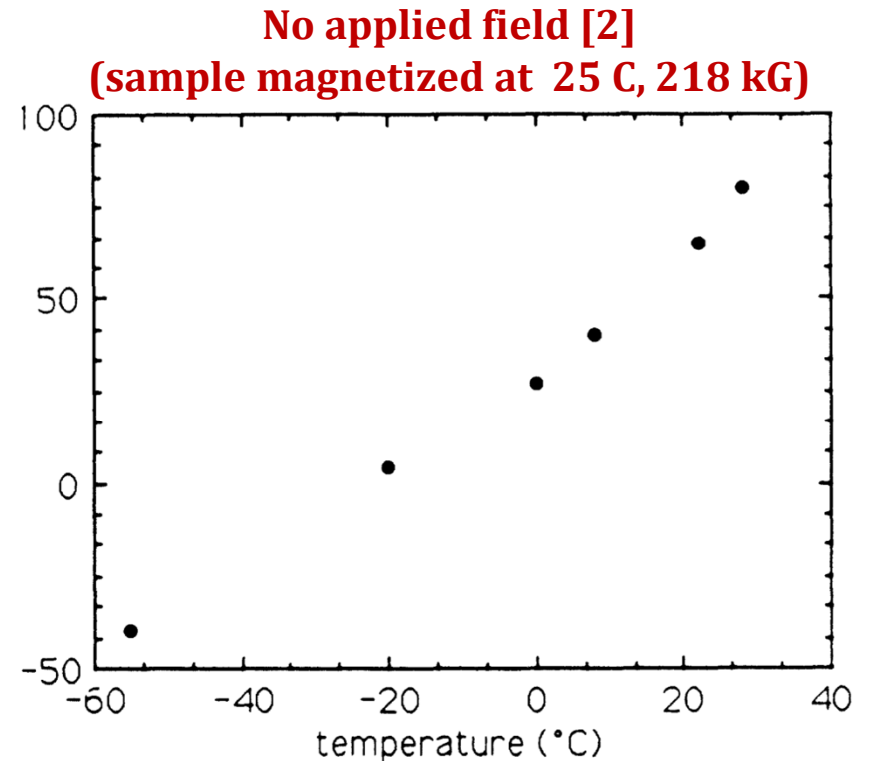


# Dy<sub>6</sub>Fe<sub>23</sub> – Magnetization vs T

- $T_c \approx 250$  K (Cool with shield, chiller, thermoelectric elements)



*Dashed: Molecular mean field theory calculation*  
*Solid: sum of individual lattice calculations*  
*Points: measurement*

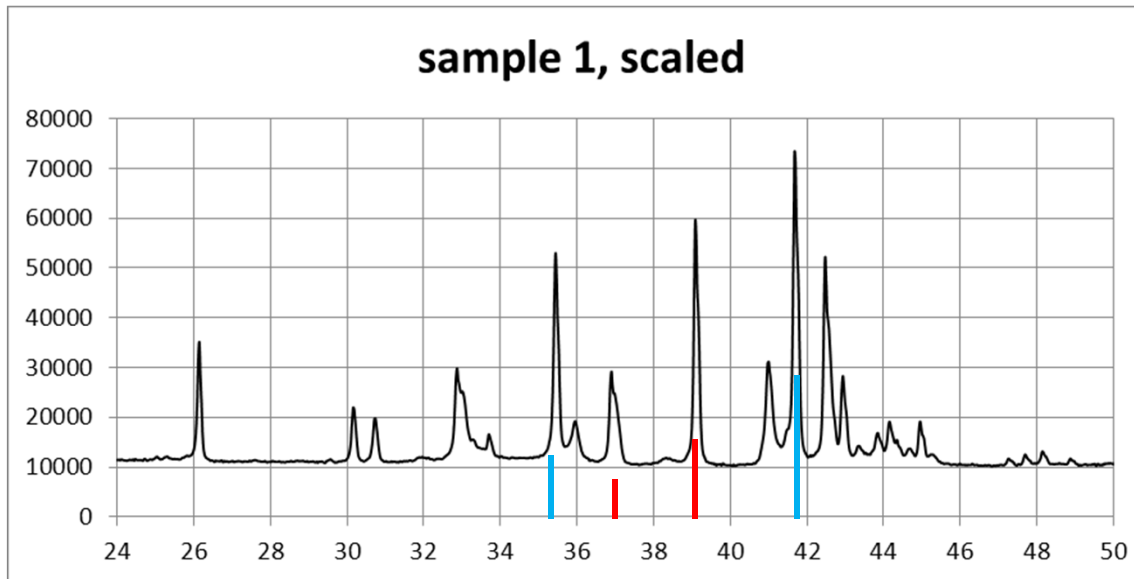


[1] J. Herbst, J. Croat, J. Appl. Phys. 55 (1984) 3023.

[2] R. Ritter, C. Goldblum, W.-T. Ni, G. Gillies, C. Speake, PRD 42 (1990) 977.

## Dy<sub>6</sub>Fe<sub>23</sub> Production (Ames National Lab)

- Melt 20.7 / 79.3 wt.% Dy/Fe in furnace
- Anneal several days at 1200°C
- X-Ray diffraction analysis:



Dy<sub>6</sub>Fe<sub>23</sub> peak

Dy<sub>1</sub>Fe<sub>3</sub> peak

Relative peak size:  
6-23 : 1-3 phase ~ 45%

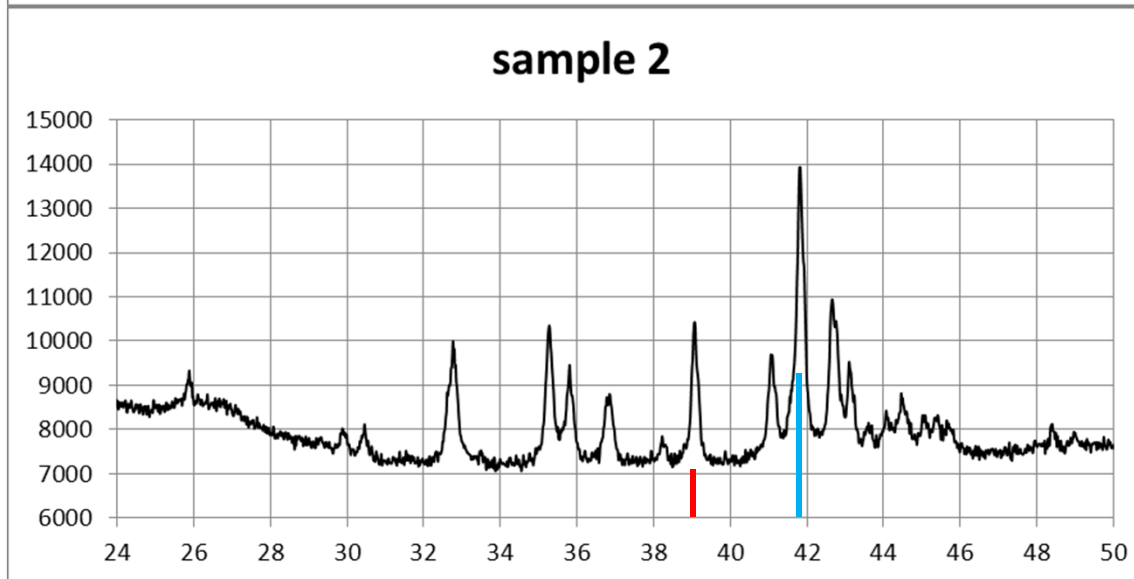
Second annealing  
(3 weeks 1200°C):

Ratio 6-23 : 1-3 reduced

- Less 6-23?
- Other phases more abundant?
- Amorphous combination?

Re-anneal in low pressure  
O<sub>2</sub> atmosphere?

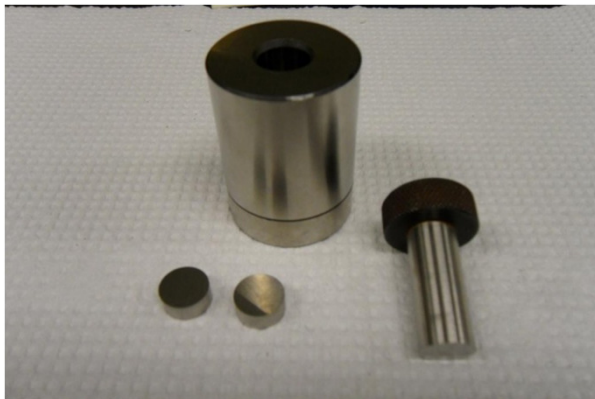
Other materials...



# $\text{Dy}_6\text{Fe}_{23}$ sample fabrication - Pressing test samples

1. **Metallic powder + binder (2g paraffin in heptane, or 10% Cereox powder)**
2. **Pour mixture into die, press hydraulically at 3000 psi**
  - **Robust sub-millimeter samples (Fe) routinely attainable**
  - **Goal: Repeat procedure with ferrimagnet**

**Die**



**Samples**



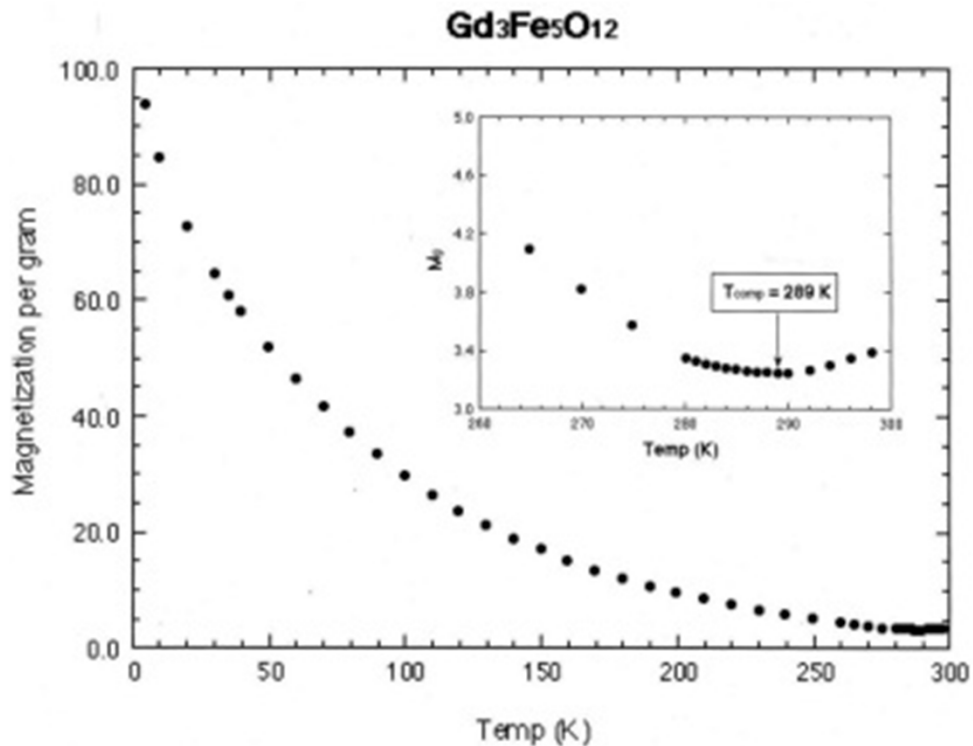
**Press**



**Next step: Magnetize in strong field, measure magnetization versus temperature**

# Rare-Earth Iron Garnets

*Geselbracht, et al., J. Chem. Edu. 71 (1994) 696*



**Combine 1 M Dy(NO<sub>3</sub>)<sub>3</sub>·6H<sub>2</sub>O with  
1 M FeCl<sub>3</sub>·6H<sub>2</sub>O**

**Add NaOH dropwise to precipitate solid**

**Decant solution, wash & dry solid**

**Press into pellet**

**Bake at 900 C in air 18-24 h**

## Summary and Outlook

**High-frequency experiment currently excludes spin-independent forces  $> 10^5$  times gravitational strength above 10 microns**

**Sensitive to forces 1000 times gravitational strength at 10 microns**

Reduce metrology uncertainties

Understand, reduce “probe-force” background

*Chameleon Analysis (2003 data)*

**Cryogenic experiment with gravitational sensitivity at 20 microns proposed**

Demonstrate cryogenic transducer and thermal noise below 10 K

**Spin-dependent experiments with same technique potentially  $\sim 8$  orders of magnitude more sensitivity than current experiments (thermal noise limit); 3-5 orders more sensitive with reasonable magnetic backgrounds**

Find a pure ferrimagnet and demonstrate  $200 < T_c < 300$  (candidates exist)

Fabricate flat samples, attach to test masses

Assemble cooling system

Demonstrate stimulated magnetic forces at  $T > T_c$

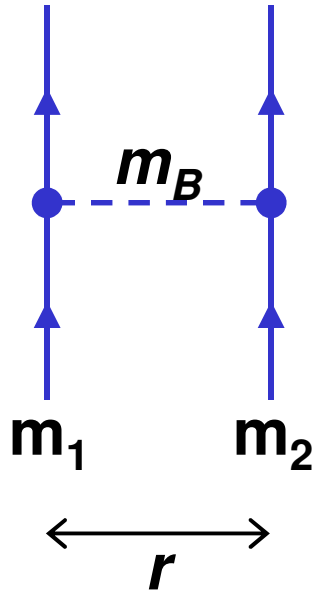
Collect data at  $T_c$

Restore thicker shield (if needed), investigate high permeability layers

**(Supplemental Slides)**

# Parameterization

## Yukawa Interaction



$$V(r) = -G \frac{m_1 m_2}{r} \left[ 1 + \alpha e^{-r/\lambda} \right]$$

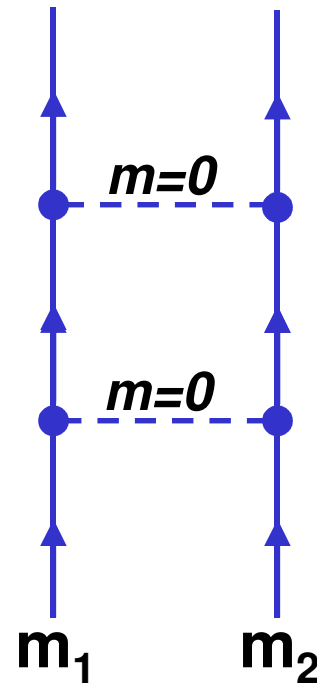
$$\lambda = \hbar / m_B c = \text{range}$$

$\alpha$  = strength relative to gravity

## Power Law

$$V(r) = -G \frac{m_1 m_2}{r} \left[ 1 + \beta_n \left( \frac{r_0}{r} \right)^{n-1} \right]$$

$r_0$  = experimental scale



set limits on  $\beta_n$  for  $n = 2 - 5$

## Signals in Recent Data

**Pre-2010: ~10 x thermal noise ( $10^3$  s), non-resonant**

- electronic pick-up, switched to differential amplifier

**Spring 2010: ~ 5x thermal noise ( $10^3$  s), resonant, position independent**

- Vibration, replaced stacks

**Summer/fall 2010: ~ 2-5 x thermal noise, resonant, weak gap dependence**

- “long-range” capacitive coupling, re-designed shield

**Spring/summer 2011: ~ 2-5 x thermal noise, resonant, smallest gaps**

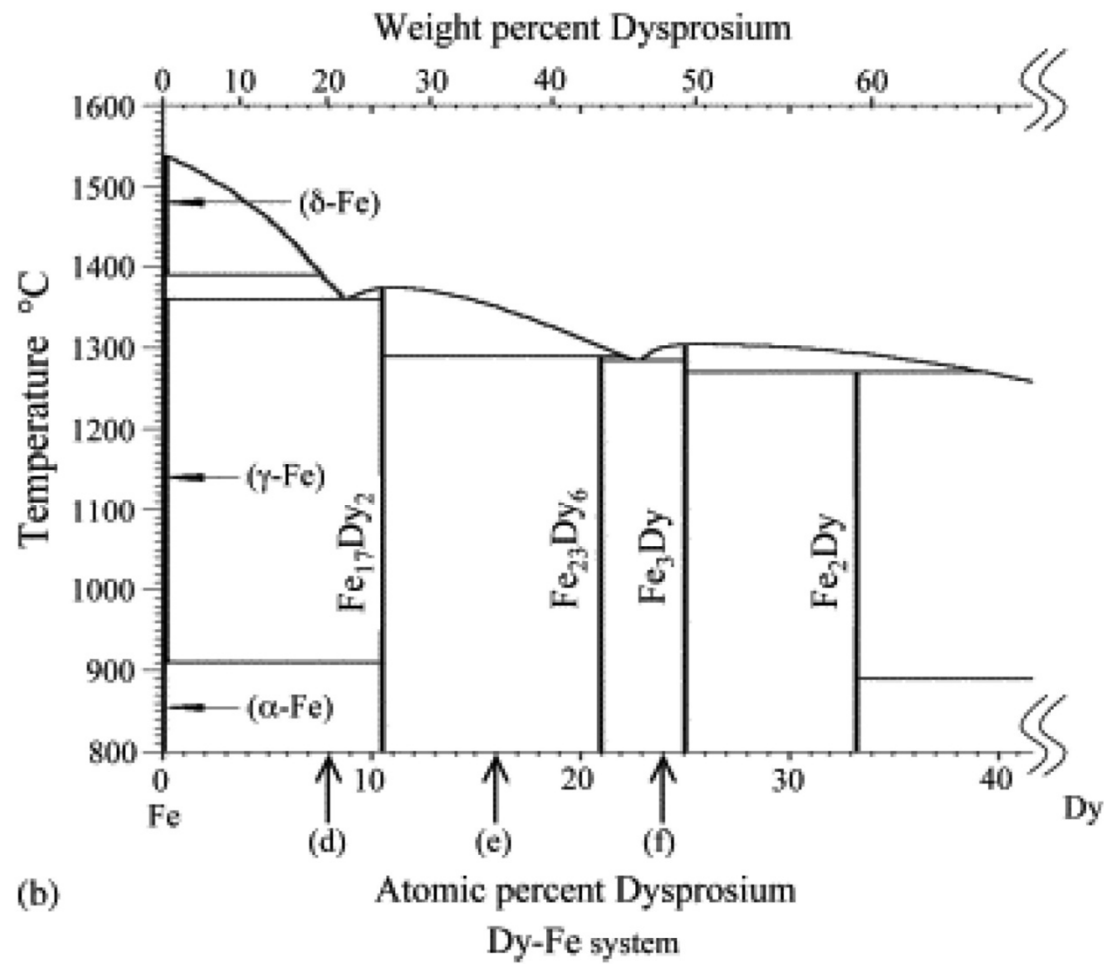
- Poorly-grounded shield or the problem below (?)

**Fall 2011: fluctuating resonant signals *and* noise**

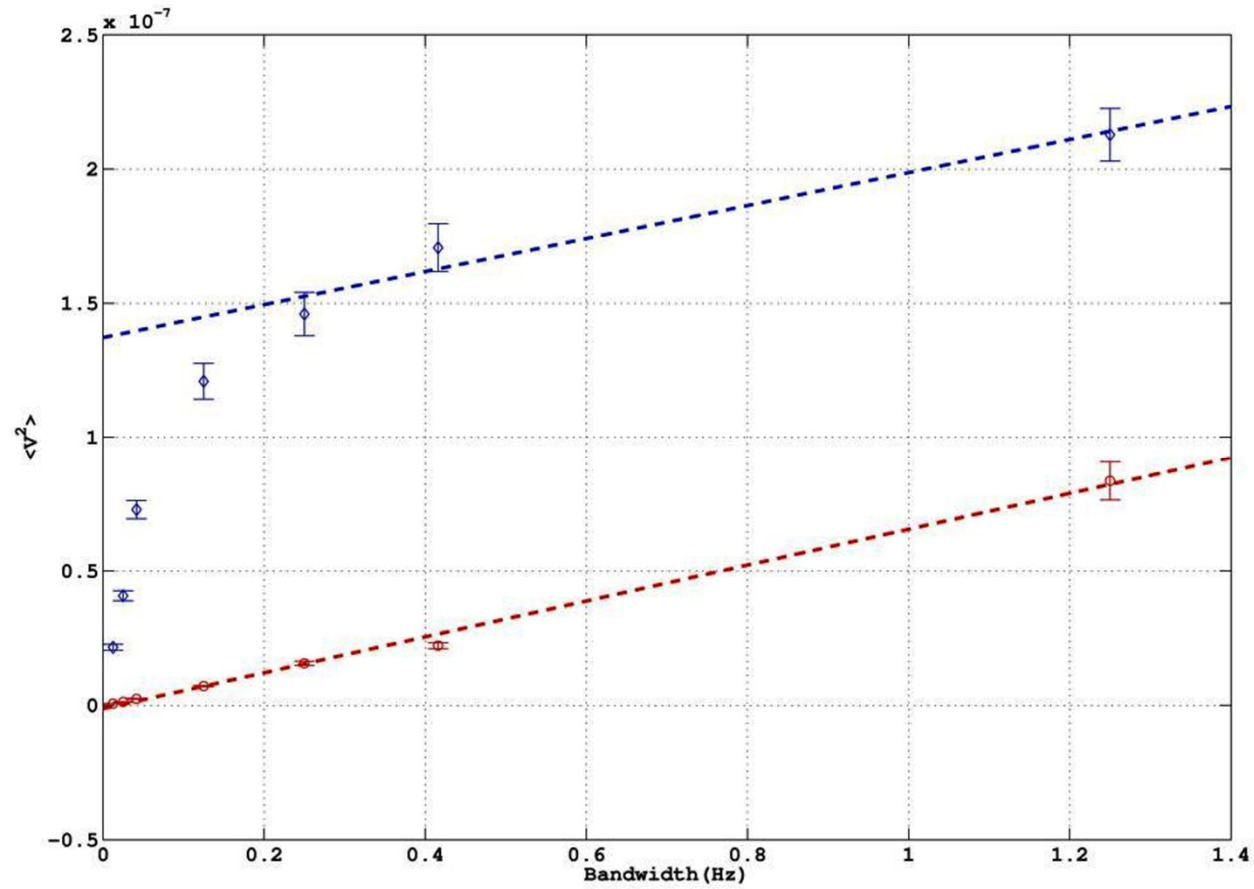
- Faulty (?) differential input on lock-in



# Dy-Fe Phase Diagram



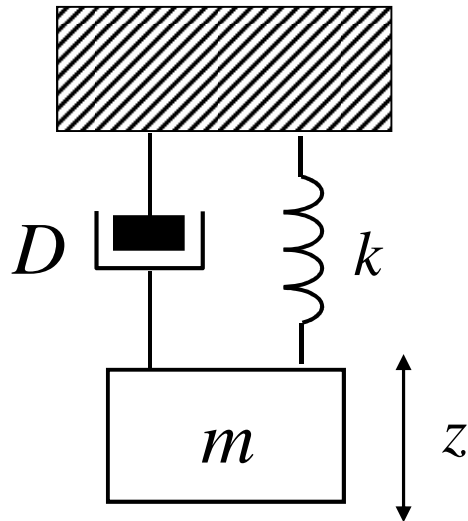
# Thermal and Amplifier Noise



**$G = 20k \rightarrow \text{amp} \sim 10 \text{ nV}/\sqrt{\text{Hz}}$**

# Calibration with Thermal Noise

## Detector Model:



## Free thermal oscillations:

$$\frac{1}{2} k_B T = \frac{1}{2} m \omega^2 z_{T(rms)}^2$$

## Damped, driven oscillations on resonance:

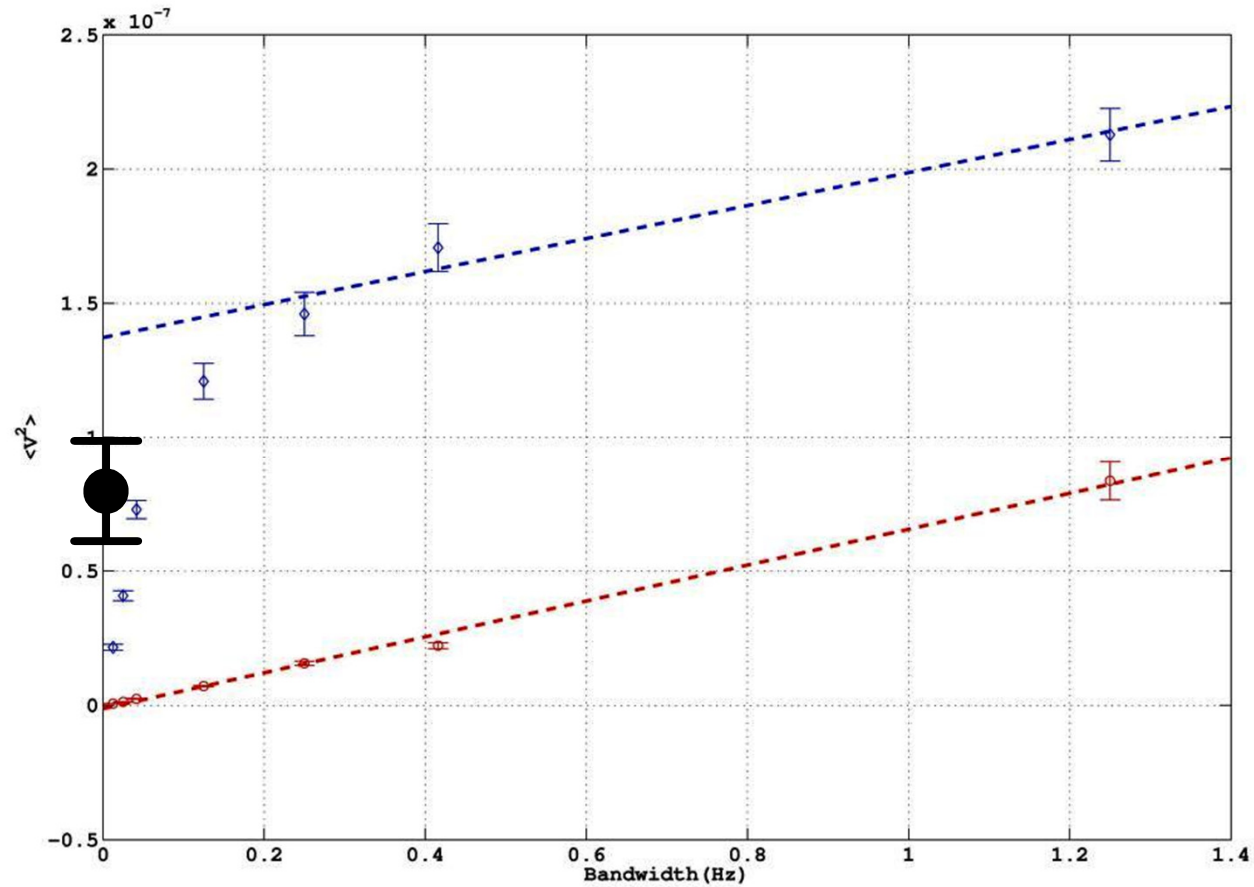
$$F_D = -\frac{m\omega^2}{Q} z_D \quad \text{where} \quad Q = \frac{m\omega}{D}$$

$$\Rightarrow \text{Measured force: } F_D = -\frac{\overline{z_D}}{z_{T(rms)}} \frac{\omega \sqrt{mk_B T}}{Q}$$

$$z_T, z_D, \omega, T, Q \text{ from data, } \frac{\overline{z_D}}{z_{T(rms)}} = \frac{\overline{V_D}}{V_{T(rms)}}$$

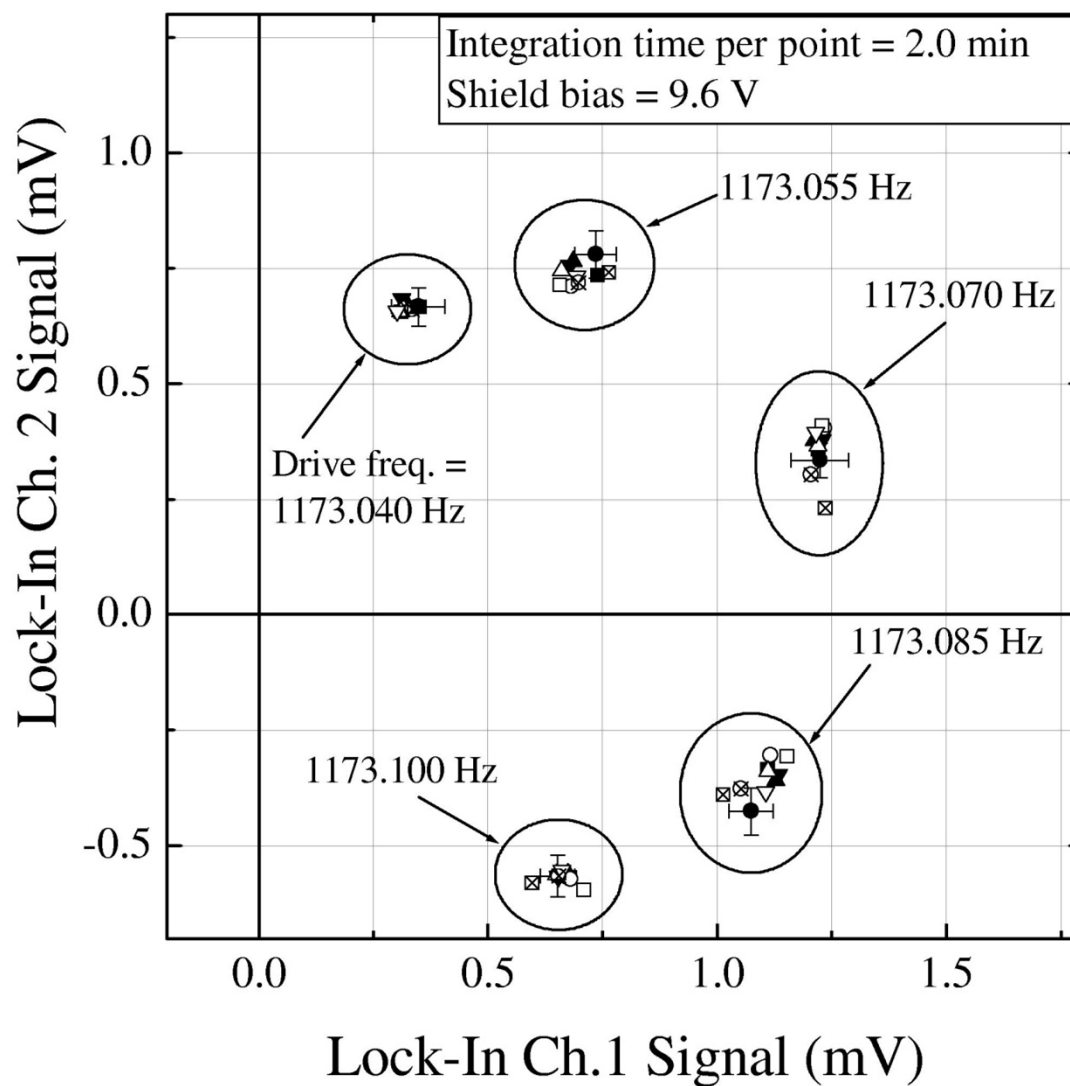
For distributed oscillator sampled at  $r$ ,  $m \rightarrow \frac{\rho \int |z|^2 dV}{|z(r)|^2}$  ← mode shape from computer model

# Thermal and Amplifier Noise



$G = 20k \rightarrow \text{amp} \sim 10 \text{ nV}/\sqrt{\text{Hz}}$

# Diagnostic Data – Shield Biased



## Analysis steps

Likelihood function:

$$L(x | \alpha, \nu) = \left( \frac{1}{\sigma \sqrt{2\pi e}} \right)^N \exp \left[ \frac{-(\bar{x} - \mu(\alpha, \nu))^2}{2(\sigma / \sqrt{N})^2} \right]$$

$\bar{x}$  = average voltage measured,  $\sigma$  = standard deviation,  $N$  = number of samples

$\mu(\alpha, \nu)$  = predicted voltage for given  $\alpha$  and set of systematics  $\nu$

General expression for predicted average voltage:

$$\mu(\alpha, \nu) = \sqrt{|V^T|^2} \frac{Q}{\omega_0 \sqrt{k_B T \rho_d}} \frac{\int d^3 \vec{r}' \vec{z}^F(\vec{r}') \cdot \vec{f}(\vec{r}')}{\sqrt{\int d^3 \vec{r}' |\vec{z}^F(\vec{r}')|^2}}$$

$V^T$  = thermal noise voltage fluctuations

$Q$  = detector mechanical quality factor

$\omega_0$  = detector resonant frequency

$\rho_d$  = detector mass density

$z^F(r')$  = displacement of detector at arbitrary point in detector  $r'$

$f(r')$  = force/unit volume on detector at arbitrary point in detector  $r'$

## Analysis steps

Interval  $[\alpha_{lo}, \alpha_{up}]$  that contains true  $\alpha$  with probability CL:

$$CL = \int_{\alpha_{lo}}^{\alpha_{up}} p(\alpha | x) d\alpha$$

$p(\alpha/x)$  = probability density function for  $\alpha$  given DATA  $x$

(For any  $\lambda$ , find  $[\alpha_{lo}, \alpha_{up}]$  so that CL = 0.95)

e.g., for  $\lambda = 20 \mu\text{m}$ ,  $-6 \times 10^3 < \alpha < 4 \times 10^3$

Bayes' Theorem:

$$p(\alpha | x) = \frac{L(x | \alpha)\pi(\alpha)}{\int_{-\infty}^{\infty} L(x | \alpha')\pi(\alpha')d\alpha'}$$

$L(\alpha/x)$  = likelihood function INTEGRATED OVER SYSTEMATICS  $v$

$\pi(\alpha)$  = prior pdf for  $\alpha$  (assumed uniform between old limits on  $\alpha$ )

# Analysis steps

Monte Carlo program calculates:

$$\int d^3\vec{r}' \vec{z}^F(\vec{r}') \bullet \vec{f}(\vec{r}')$$

$z^F(r')$  = displacement of detector at arbitrary point  $r'$  in detector

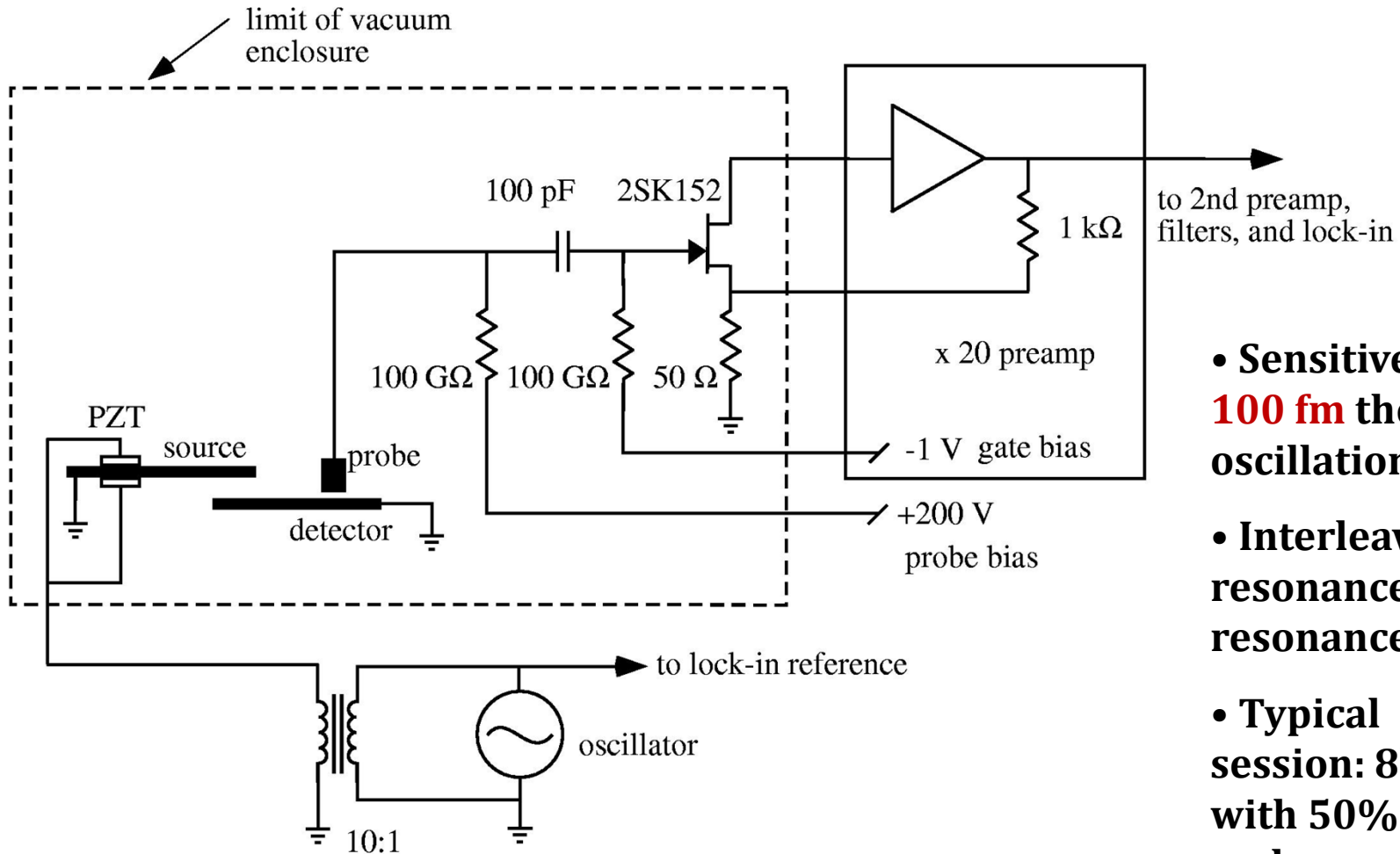
(given by 2<sup>nd</sup> order polynomial fit to geometry survey data)

$f(r')$  = force/unit volume on detector at arbitrary point  $r'$  in detector

What is  $f(r')$  due to interaction with source mass for case of LLV force?

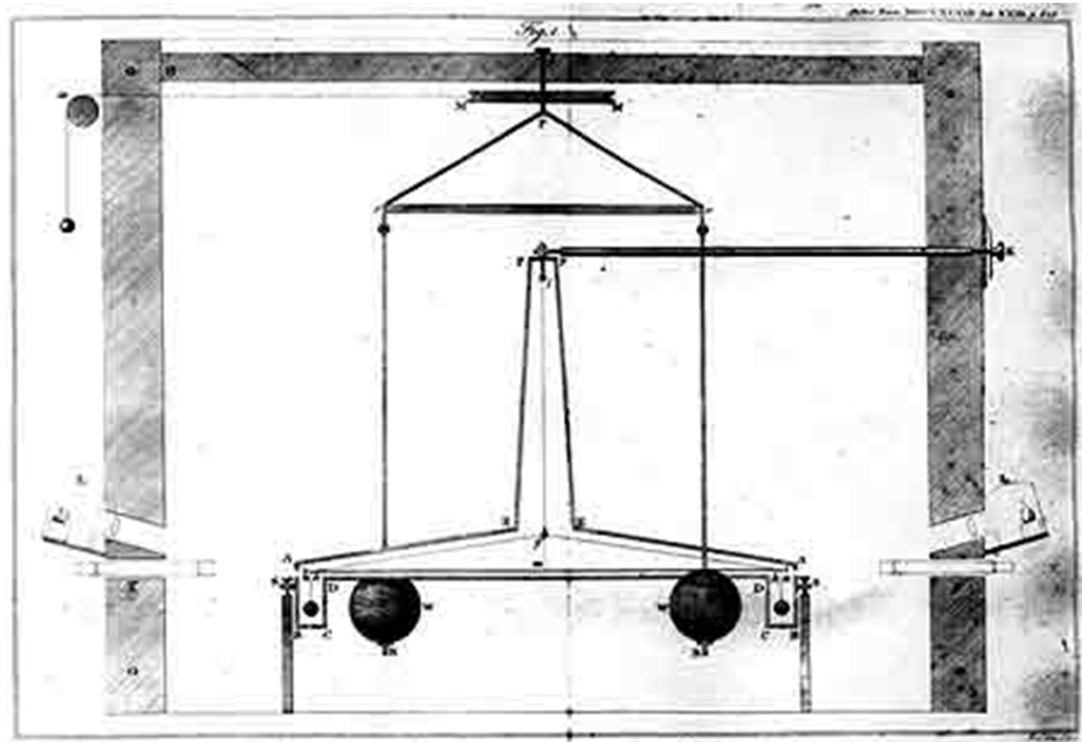
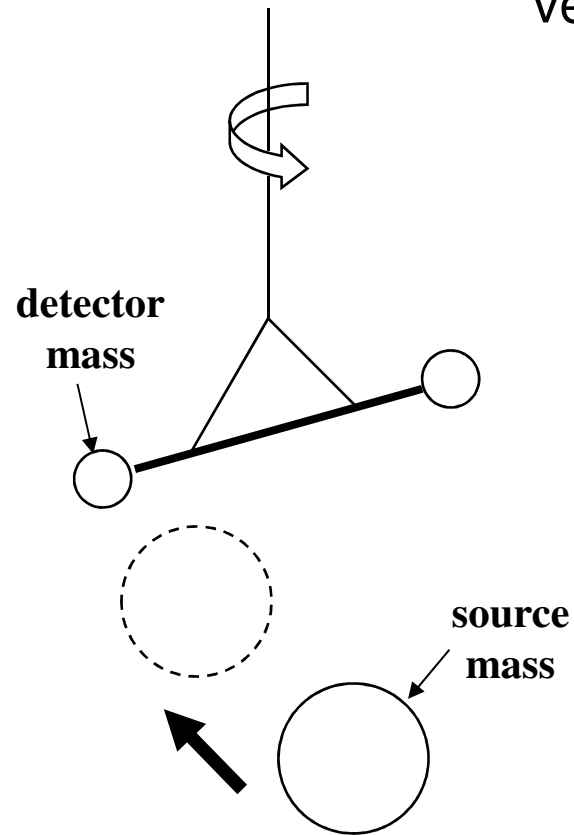


# Readout



- Sensitive to  $\approx$  **100 fm** thermal oscillations
- Interleave on resonance, off resonance runs
- Typical session: 8hrs with 50% duty cycle

Experiment is short-range ( $\sim 50 \mu\text{m}$ ), high-frequency (1 kHz)  
version of Cavendish Experiment

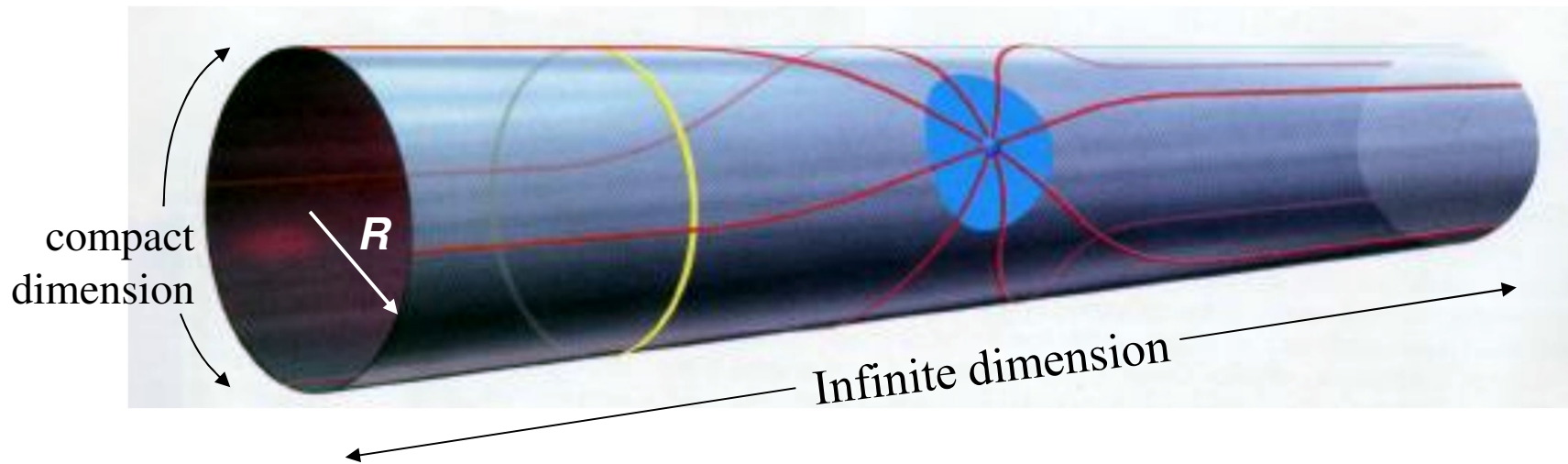


Pb test masses ( $\rho = 11000 \text{ kg/m}^3$ ): large = 20 cm diam., small = 5 cm diam

$$G = 6.76 \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$$

2006 CODATA:  $G = 6.67428(67) \times 10^{-11} \text{ m}^3 \text{ kg}^{-1} \text{ s}^{-2}$

# “Large” Extra Dimensions



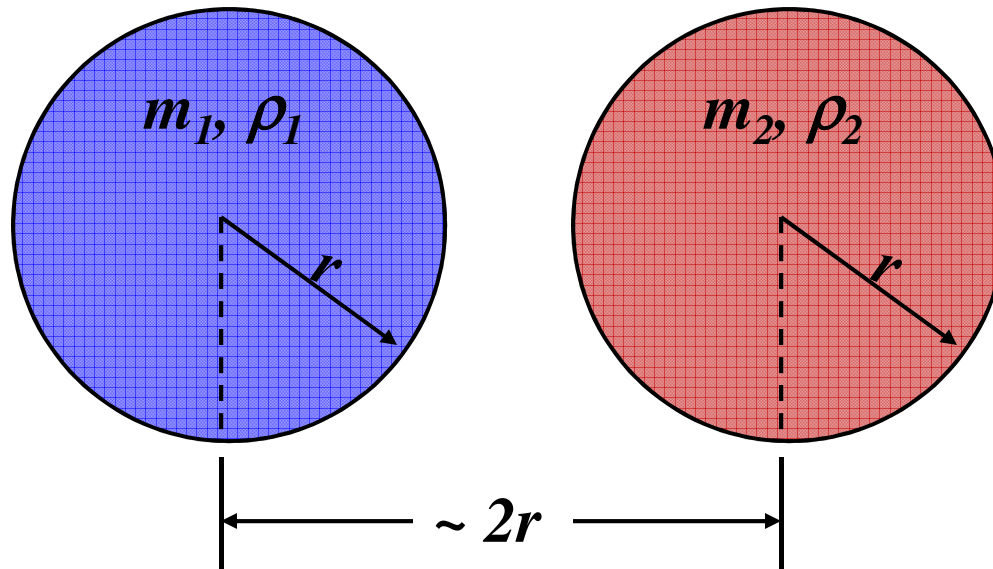
**Strong, Weak, EM force confined to 3 dimensions**

- **Gravity spreads out into  $n$  extra dimensions of size  $R$ , appears diluted**

$$R = \left[ \frac{M_P}{M^*} \right]^{2/n} \left[ \frac{\hbar}{2\pi M^*} \right]$$

- **Gravity unifies with EW force ( $M^* \sim 1$  TeV) if  $n = 2$ ,  $R \sim 1$  mm  
 $n = 3$ ,  $R \sim 1$  nm**

## Challenge: Scaling with Size of Apparatus



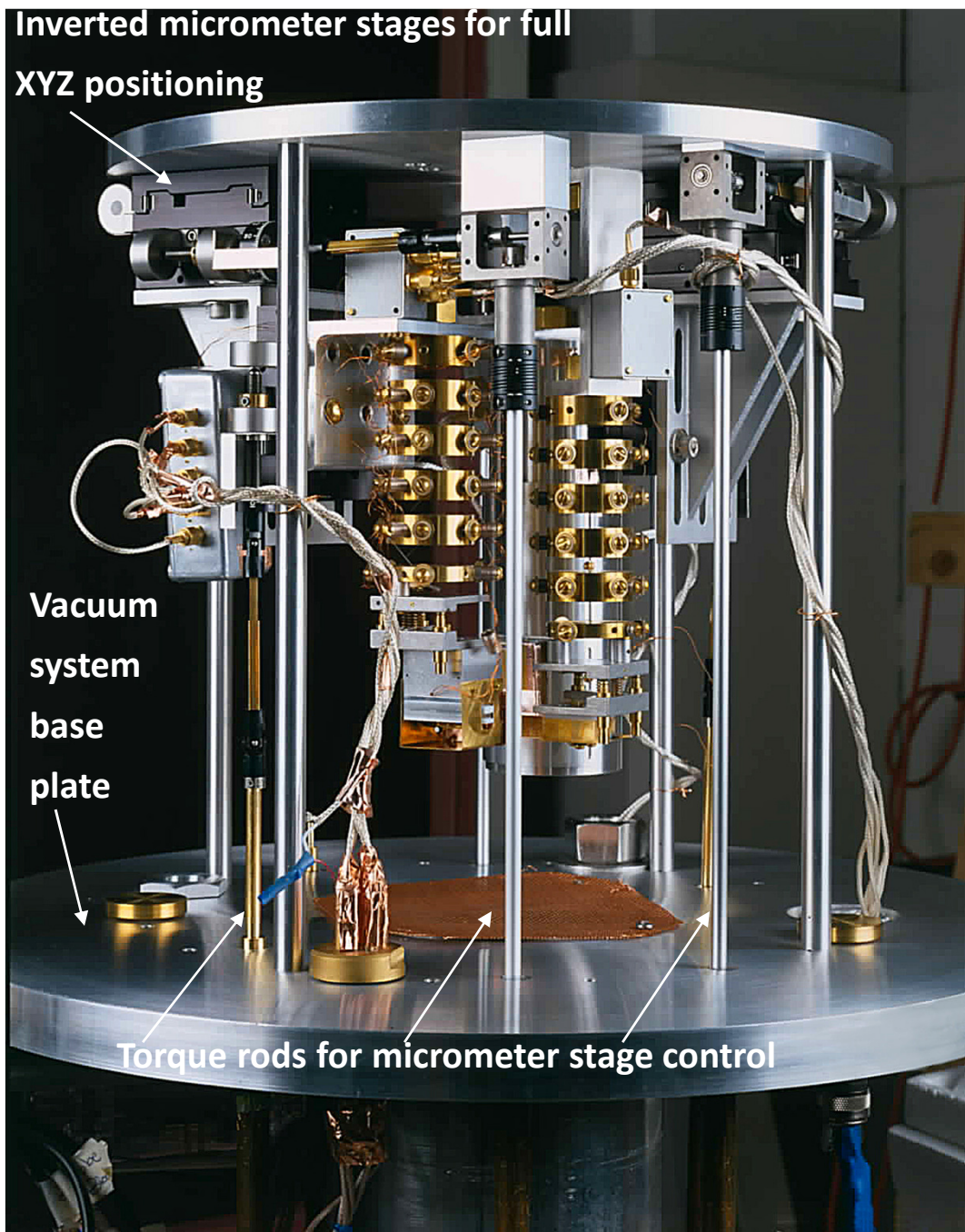
$$F = \frac{Gm_1m_2}{(2r)^2} = \frac{G\rho_1\rho_2(4r^3)^2}{4r^2} \sim G\rho_1\rho_2r^4$$

$$\rho_1 = \rho_2 = 20 \text{ g/cm}^3, r = 10 \text{ cm} \Rightarrow F \approx 10^{-5} \text{ N}$$

$$r = 100 \text{ }\mu\text{m} \Rightarrow F \approx 10^{-17} \text{ N}$$

**Background Forces:**  $\sim r^{-2}$  (electrostatics),  $\sim r^{-4}$  (magnetic dipoles, Casimir)

# Central Apparatus



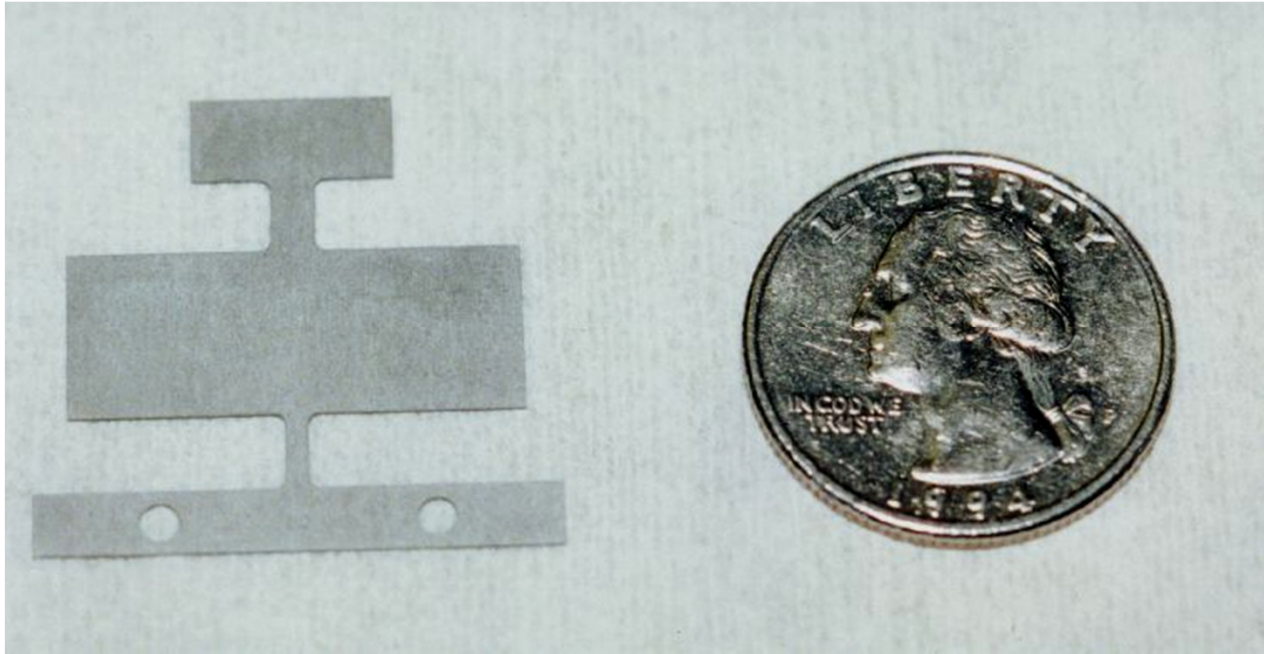
# Installation at IUCEEM

## Vacuum System

- Hollow riser for magnetic isolation
- LN<sub>2</sub> - trapped diffusion pump mounts below plate
- $P \sim 10^{-7}$  torr



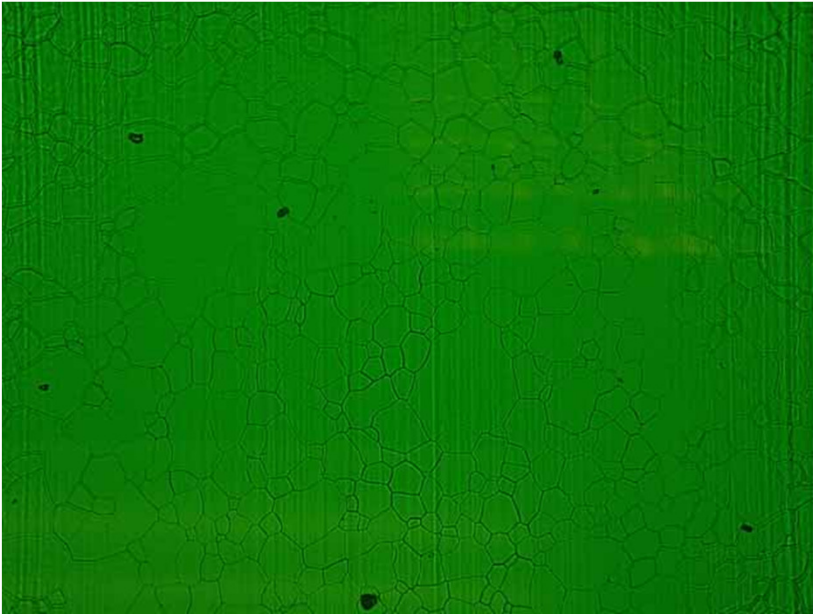
New detector prototypes have been fabricated



- **200  $\mu\text{m}$  thick tungsten sheet (high density)**
- **Fabricated by wire EDM**
- **First generation: Annealed at 1600 K in helium atmosphere**
- **New oscillators: Annealed at 2700 K; expect larger crystals, higher Q**

2700 K annealing leads to much larger crystals

**New detector surface  
(200 x magnification)**



**1000 x showing 90  $\mu\text{m}$  crystal  
(previous maximum = 15  $\mu\text{m}$ )**



**Higher T anneals had expected material effect, but  
mechanical properties still under test...**



# New detectors and projected sensitivity

## Available prototypes

| <i>Material</i>    | <i>Q @ 300 K</i>  | <i>Q @ 77 K</i>   | <i>Q @ 4 K</i>       |
|--------------------|-------------------|-------------------|----------------------|
| Si                 | $6 \times 10^3$   | $1 \times 10^5$   | $8 \times 10^5$      |
| W (as machined)    | $7 \times 10^3$   | $1 \times 10^3$   | $1.2 \times 10^4$    |
| W (1600 K anneal)* | $2.5 \times 10^4$ |                   |                      |
| W (2700 K anneal)  | $2.8 \times 10^4$ | $1.8 \times 10^5$ | $1 \times 10^6$ (8K) |

*\*Used for published experiment*

Data of W. Duffy (~ 3 cm diameter, 1 kHz cylindrical torsional oscillators):

*J. Appl. Phys. 72 (1992) 5628*

| <i>Material</i>   | <i>Q @ 300 K</i> | <i>Q @ 77 K</i> | <i>Q @ 4 K</i>  |
|-------------------|------------------|-----------------|-----------------|
| W (as machined)   | $2 \times 10^4$  | $1 \times 10^5$ | $5 \times 10^5$ |
| W (2023 K anneal) | $2 \times 10^5$  | $1 \times 10^6$ | $1 \times 10^7$ |

$\alpha \sim \sqrt{T/Q}$  improves by factor of 50 at 4 K

# New detectors and projected sensitivity

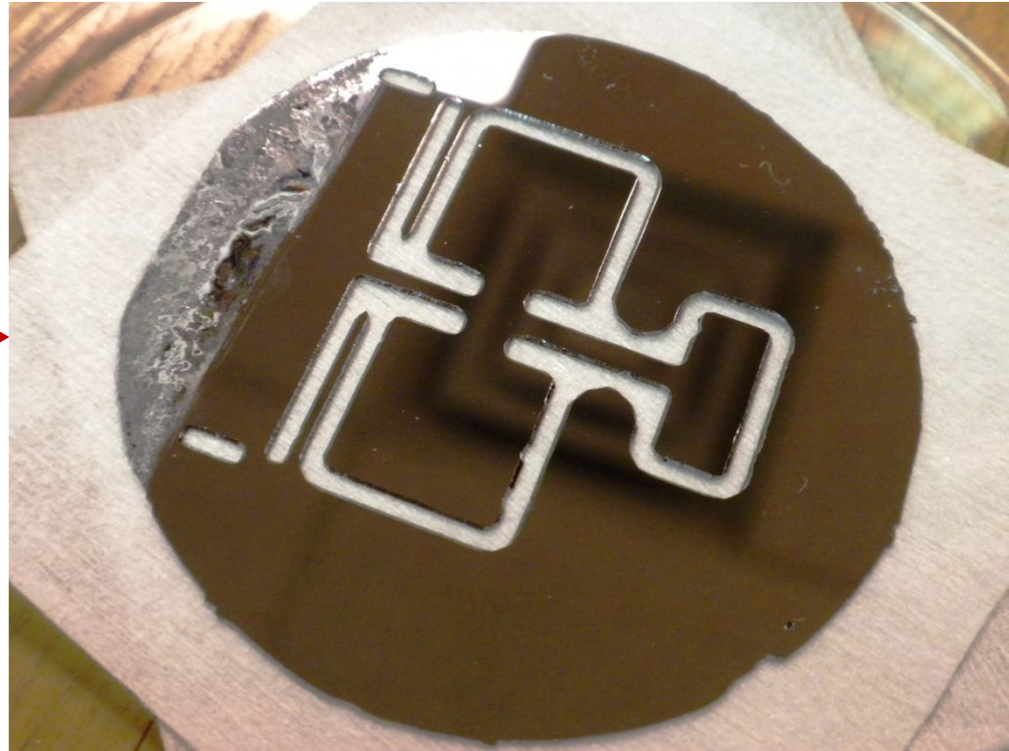
## Thinner test masses

Reduce Newtonian background

Solid, 30  $\mu\text{m}$  W too flimsy

200  $\mu\text{m}$  Si with 30  $\mu\text{m}$  gold plate

S. Jacobson, IU Chemistry

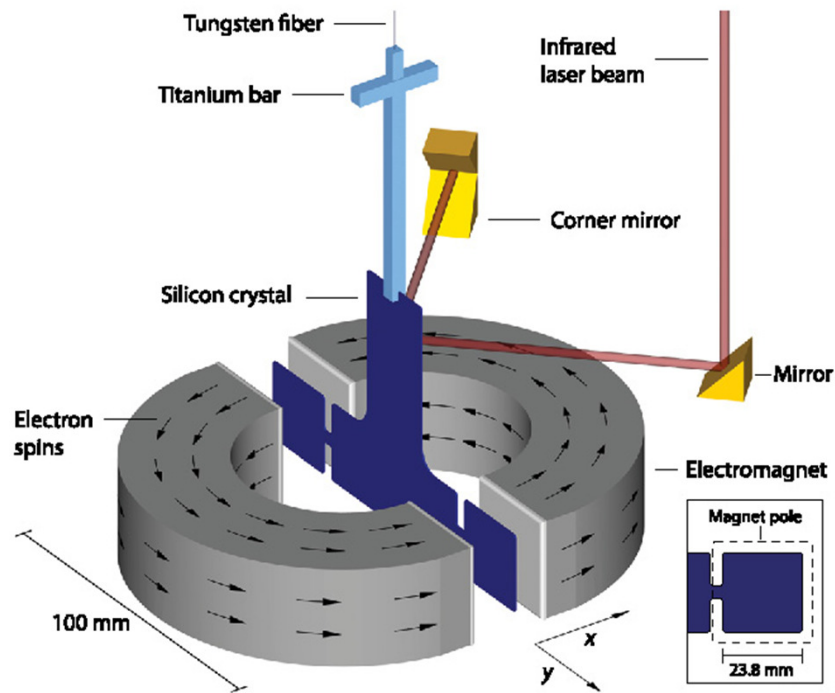


## Available prototypes

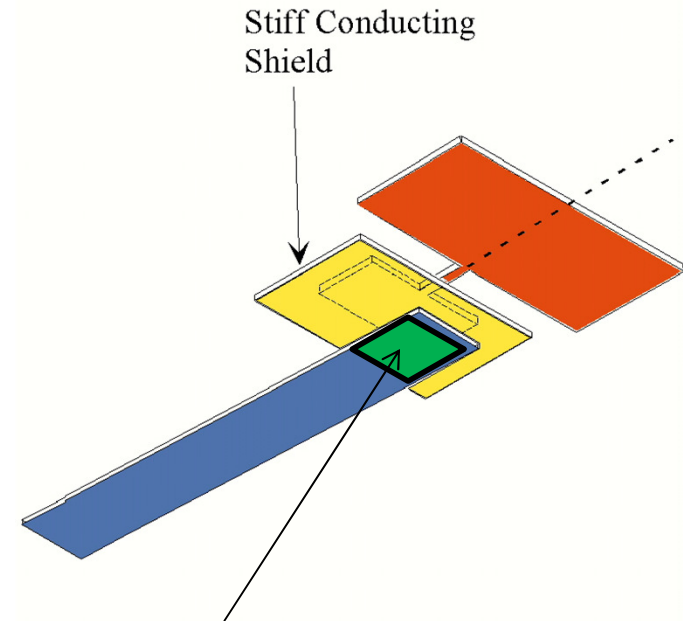
| <i>Material</i>   | <i>Q @ 300 K</i>  | <i>Q @ 77 K</i>   | <i>Q @ 4 K</i>       |
|-------------------|-------------------|-------------------|----------------------|
| Si                | $6 \times 10^3$   | $1 \times 10^5$   | $8 \times 10^5$      |
| W (2700 K anneal) | $2.8 \times 10^4$ | $1.8 \times 10^5$ | $1 \times 10^6$ (8K) |
| Si (new)          | $1.0 \times 10^4$ | ?                 | ?                    |

# Spin – Dependent Interactions

## Eot-Wash ALP torsion pendulum



*S. Hoedl et al., PRL 106 (2011) 041801*



- **Compensated test mass ( $Dy_6Fe_{23}$ )**

**Assume 10% of attained spin densities**

**1 mm thick**

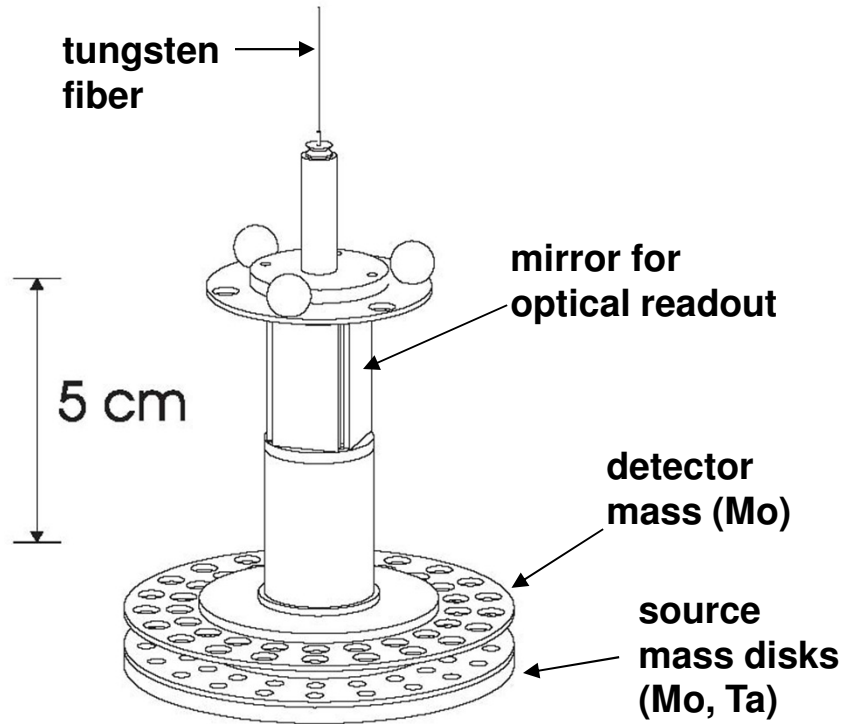
- **Magnetic shield**

**$\mu$ -metal?**

**~ 100  $\mu$ m thick**

# Eot-Wash Torsion Pendulum Experiment

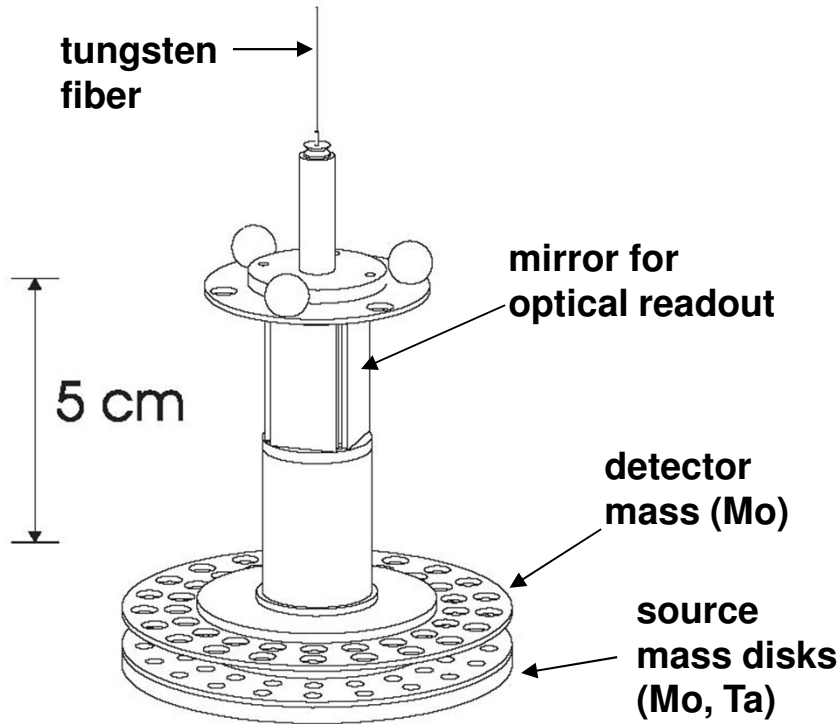
*D. Kapner, E. Adelberger et al., PRL 98 021101 (2007)*



- 55  $\mu\text{m}$  minimum gap
- 10  $\mu\text{m}$  BeCu membrane (not shown)

# Eot-Wash Torsion Pendulum Experiment

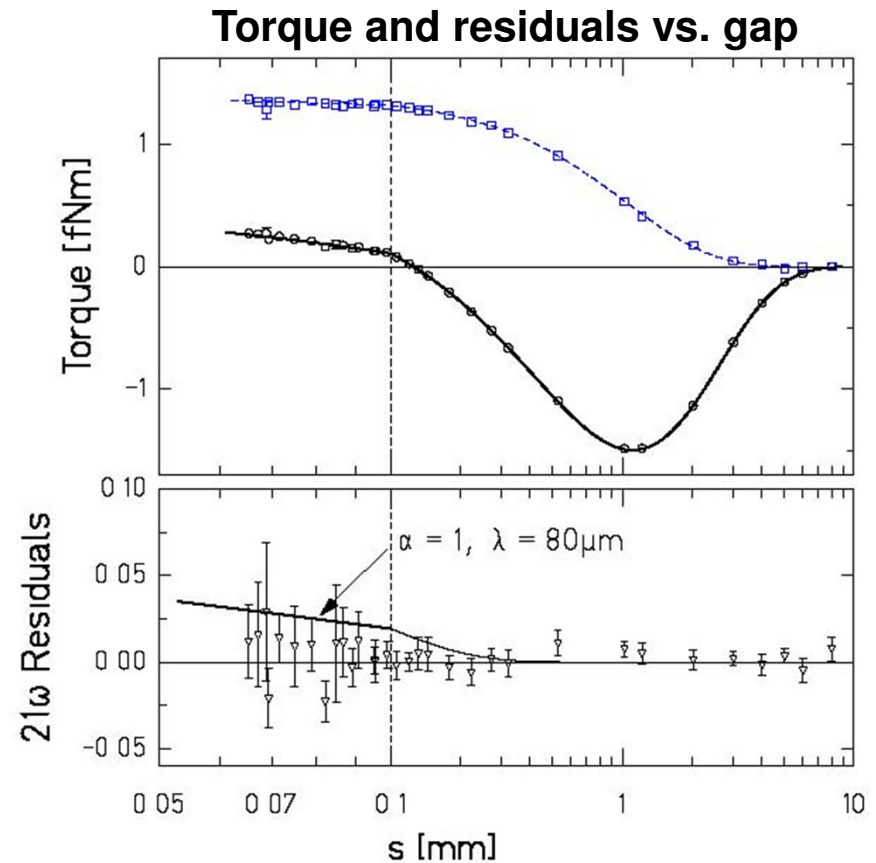
*D. Kapner, E. Adelberger et al., PRL 98 021101 (2007)*



- 55  $\mu\text{m}$  minimum gap
- 10  $\mu\text{m}$  BeCu membrane (not shown)

**Limits: Scenarios with  $\alpha \geq 1$  excluded at 95% CL for  $\lambda \geq 56 \mu\text{m}$**

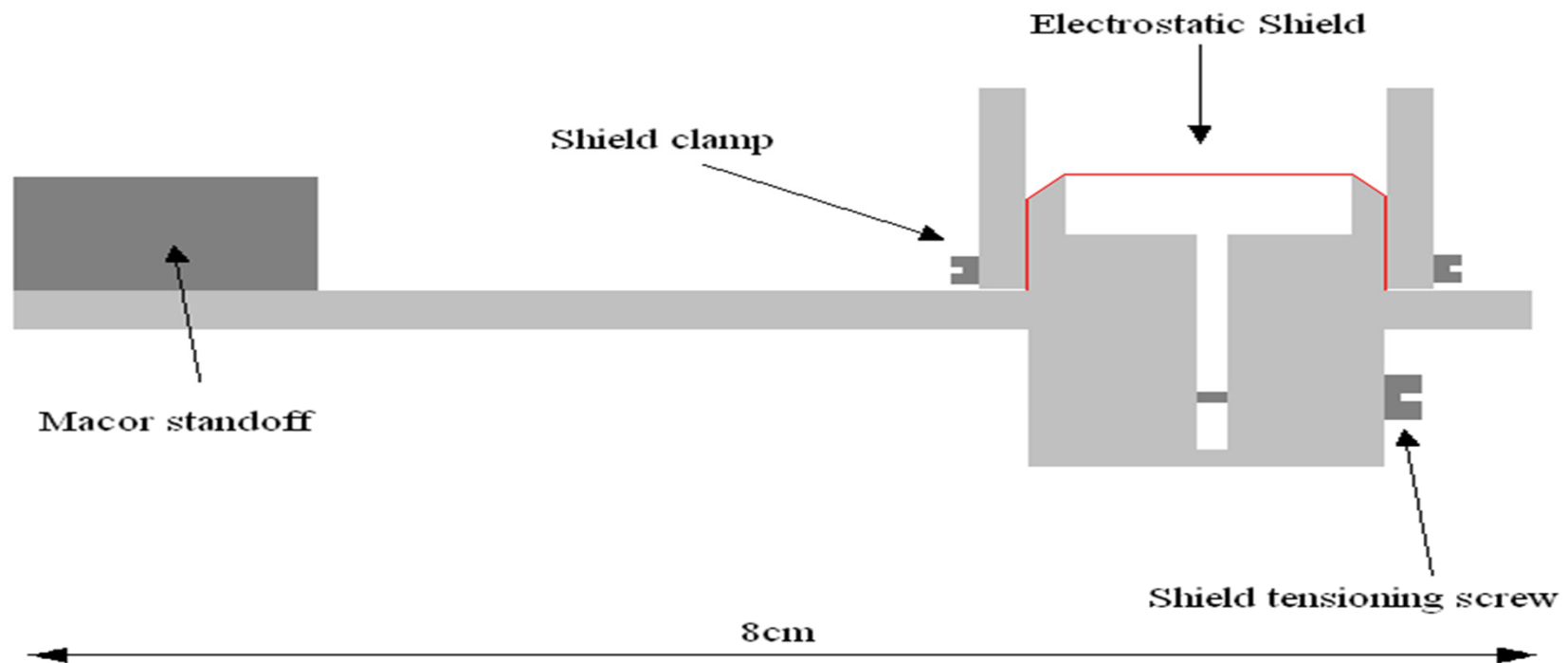
**Largest extra dimension:  $R < 44 \mu\text{m}$**



**ADD Model (2 equal-sized extra dimensions compactified on a torus):**  
 $R < 56 \mu\text{m} \Rightarrow M^* \geq 3.2 \text{ TeV}$

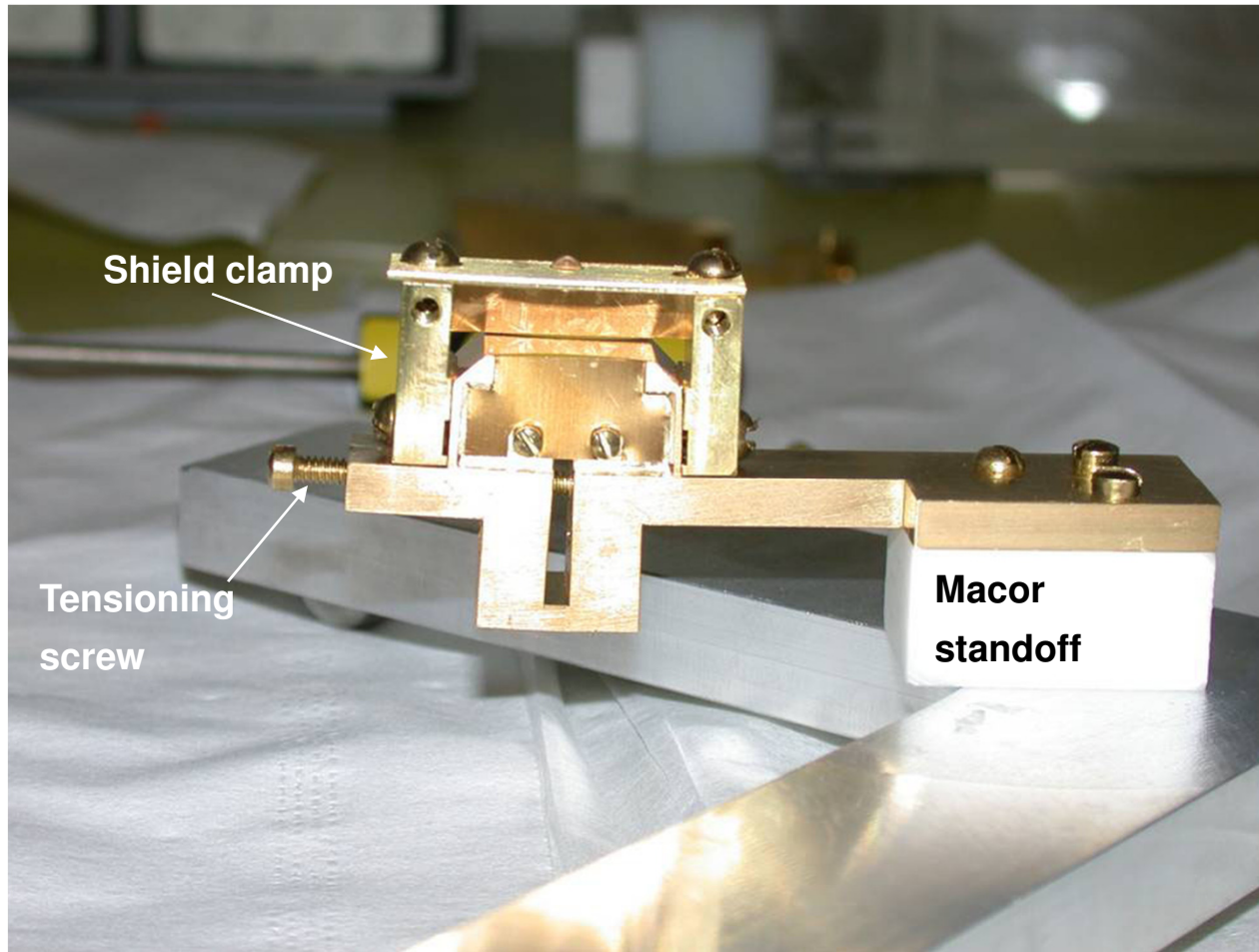
# Stretched membrane shield is designed...

- Copper-beryllium alloy stretched over frame
- Ten microns thick
- Hourglass shape for uniform distribution of tension



*Figure not to scale*

# Shield prototype exists and may be useable



- **Surface variations:**
  - 5  $\mu\text{m}$  peaks
  - 0.7  $\mu\text{m}$  rms variations (should be sufficient for  $\sim$  30  $\mu\text{m}$  experiment)
- **Conducting planes surround both test masses on 5 sides (get rid of copper tape)**
- **11/06: minimum gap = 48 microns**

**The Reality: Background forces took ~ 2 years to characterize and suppress in 1<sup>st</sup> experiment; will probably arise again at shorter ranges**

## **Vibrations**

**Filter with passive isolation stacks**

**Check that signals are geometry independent**

## **Residual Gas**

**Suppress with shield, high vacuum**

**Study with vacuum control**

## **Magnetic Forces (contaminants, eddy currents)**

**Use non-magnetic materials**

**Study with applied gradients, insulating test masses**

## **Electrostatic forces**

**Suppress with shield**

**Study with applied potentials, capacitance measurements, geometry**



# Systematic Errors

| Parameter   | Mean                        | Error                 | Units                                     |
|---|-----------------------------|-----------------------|---|
| Gravitational constant, $G$   | $6.673 \times 10^{-11}$     | $1.0 \times 10^{-13}$ | $\text{m}^3 \text{kg}^{-1} \text{s}^{-2}$ |
| Boltzmann constant, $k_B$   | $1.3806503 \times 10^{-23}$ | $2.4 \times 10^{-29}$ | $\text{J K}^{-1}$                         |
| Detector density (tungsten), $\rho_d$   | $1.93 \times 10^4$          | $1.9 \times 10^3$     | $\text{kg m}^{-3}$                        |
| Source density (tungsten), $\rho_s$   | $1.93 \times 10^4$          | $1.9 \times 10^3$     | $\text{kg m}^{-3}$                        |
| Thermal noise voltage, $\sqrt{ V^T ^2}$   | $6.09 \times 10^{-5}$       | $2.3 \times 10^{-6}$  | V   |
| Mechanical quality factor, $Q$  | $2.5522 \times 10^4$        | 29                    | (NA)                                      |
| Resonant frequency, $\omega_0/2\pi$   | 1173.085                    | 0.015                 | Hz  |
| Temperature, $T$  | 305.0                       | 0.1                   | K   |
| Integrated rms free modeshape, $\sqrt{\int d^3 \vec{r}^{\prime}  \dot{z}^F(\vec{r}^{\prime}) ^2}$ | $5.87 \times 10^{-11}$      | $5.9 \times 10^{-12}$ | $\text{m}^{5/2}$                          |

| Parameter                 | Mean ( <b>m</b> )        | Error ( <b>m</b> )   |
|---------------------------|--------------------------|----------------------|
| Detector length, $l_d$    | $5.0800 \times 10^{-3}$  | $6.4 \times 10^{-6}$ |
| Detector width, $w_d$     | $1.14550 \times 10^{-2}$ | $6.4 \times 10^{-6}$ |
| Detector thickness, $t_d$ | $1.950 \times 10^{-4}$   | $6.4 \times 10^{-6}$ |
| Source width, $w_s$       | $7.0000 \times 10^{-3}$  | $6.4 \times 10^{-6}$ |
| Source thickness, $t_s$   | $3.048 \times 10^{-4}$   | $6.4 \times 10^{-6}$ |
| Touch gap, $g_{sd}$       | $1.080 \times 10^{-4}$   | $6.4 \times 10^{-6}$ |
| Source amplitude, $dz_s$  | $1.87 \times 10^{-5}$    | $3.2 \times 10^{-6}$ |

# Consistency checks

Additional runs:

**Larger test mass gap**

**Source over opposite side of detector (and shield)**

**Reduced overlap**

- $F_{\text{pressure}} \sim F_{\text{magnetic}} \sim r^{-2}$ ,  $F_{\text{electrostatic}} \sim r^{-4}$ ,  $F_{\text{vibrational}} \sim$  (constant)
- **Shield response**

*No resonant signal observed*

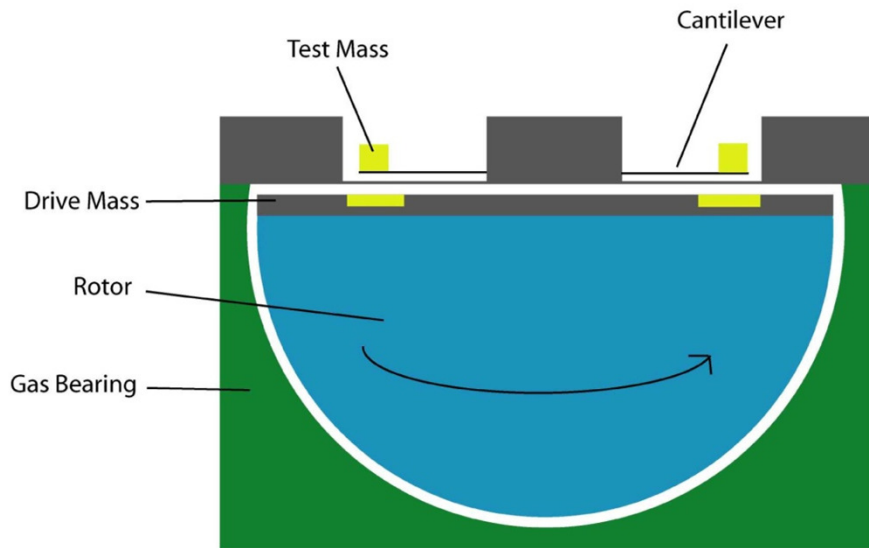
Expected backgrounds from ambient fields:

**Magnetic Background = Signal with applied  $B \times (B_{\text{ambient}} / B_{\text{applied}})^2 = 10^{-7} \text{ V}$**

**ES Background = Signal with applied  $V \times (V_{\text{ambient}} / V_{\text{applied}})^4 = 10^{-10} \text{ V}$**

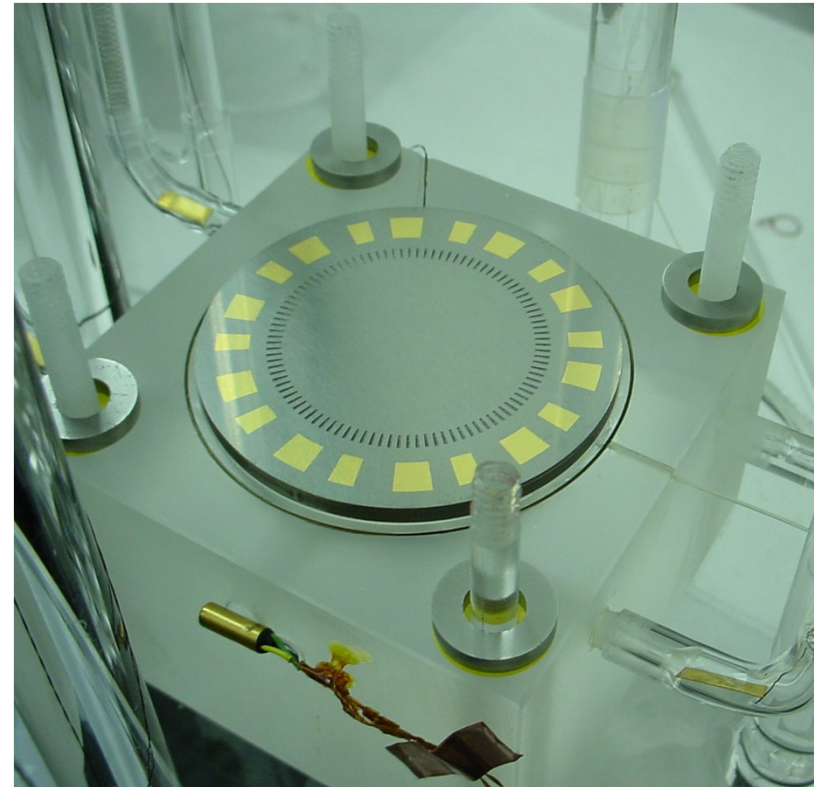
*All < thermal noise ( $10^{-6} \text{ V}$ )*

# Stanford Microcantilever Experiment – Generation II



- Masses modulated on spinning rotor
- Larger area drive and test masses for increased sensitivity

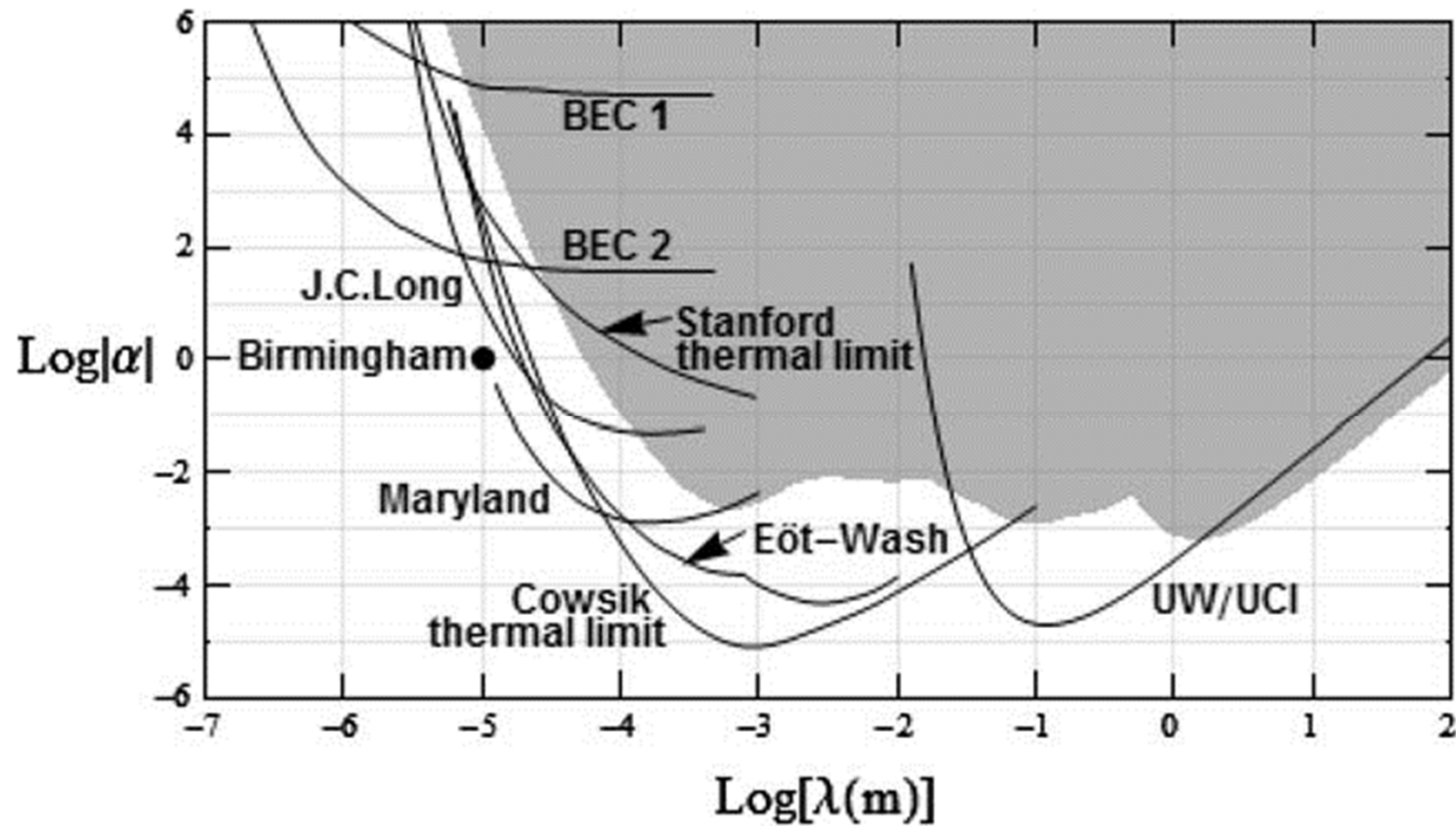
**Drive mass mounted in gas bearing**



*figures courtesy of David M. Weld*

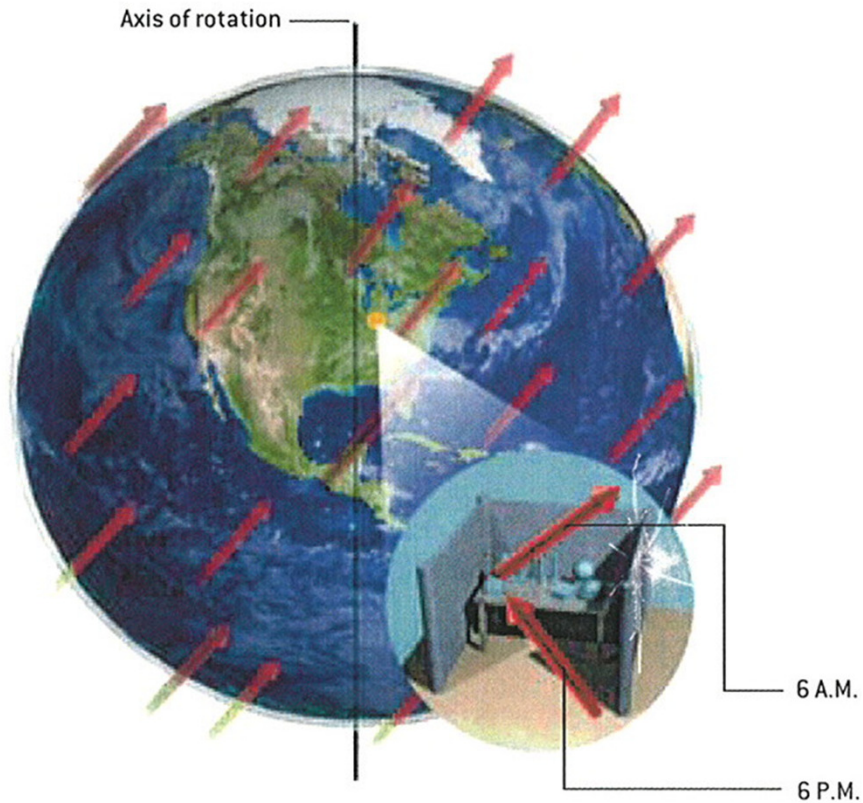
# Limits and Projections – 1 $\mu\text{m}$ – 1 m

*R. Newman, Space Sci. Rev. 148 (2009) 175*



# Search for Lorentz Violation

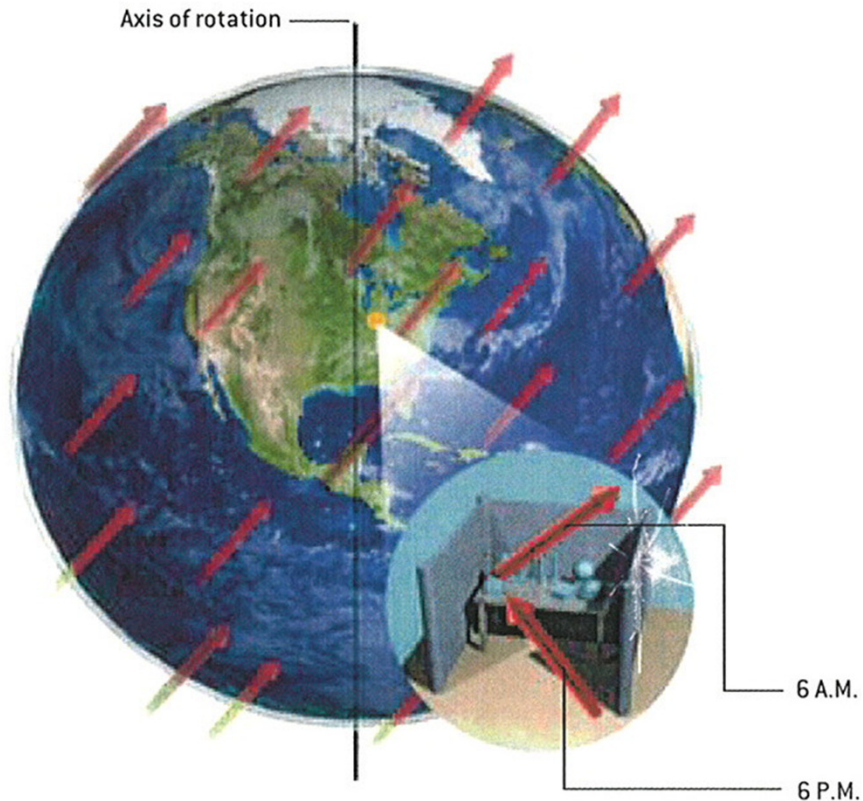
Test for sidereal variation in force signal:



Source: A. Kostelecký, *Scientific American*,  
September 2004, 93.

# Search for Lorentz Violation

Test for sidereal variation in force signal:



Source: A. Kostelecký, *Scientific American*, September 2004, 93.

Standard Model Extension (SME)

Recently expanded to gravitational sector

V. A. Kostelecký, *PRD* 69 105009 (2004).

Q. G. Bailey and V. A. Kostelecký, *PRD* 74 045001 (2006).

Action:

$$S = S_{GR} + S_{LV} + S_{MATTER}$$

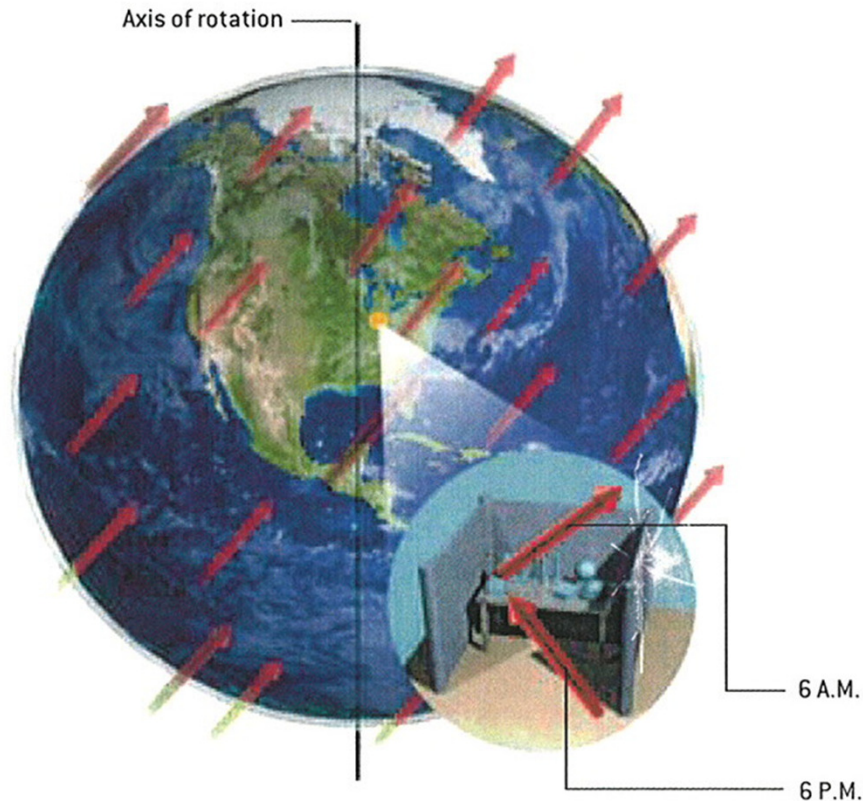
$$S_{LV} = f(\underbrace{u, s^{\mu\nu}, t^{\kappa\lambda\mu\nu}})$$

20 coefficients controlling L.V.

Estimated sensitivities:  $10^{-15} - 10^{-4}$

# New Analysis - Search for Lorentz Violation (2002 Data)

Test for sidereal variation in force signal:



Source: A. Kostelecký, *Scientific American*, September 2004, 93.

Standard Model Extension (SME)

Recently expanded to gravitational sector

*V. A. Kostelecký, PRD 69 105009 (2004).*

*Q. G. Bailey and V. A. Kostelecký, PRD 74 045001 (2006).*

Action:

$$S = S_{GR} + S_{LV} + S_{MATTER}$$

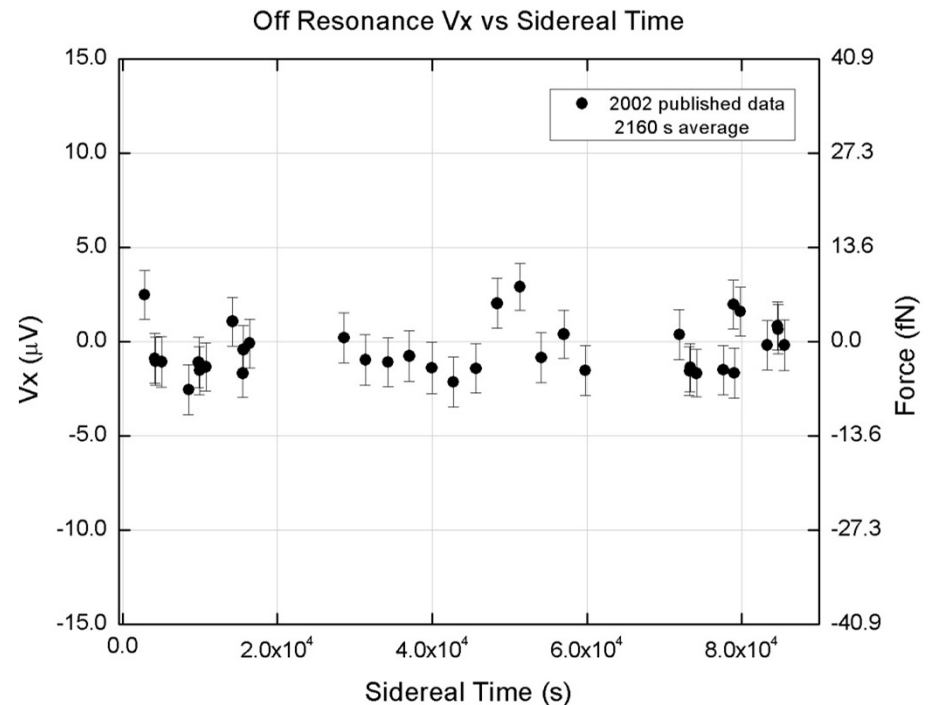
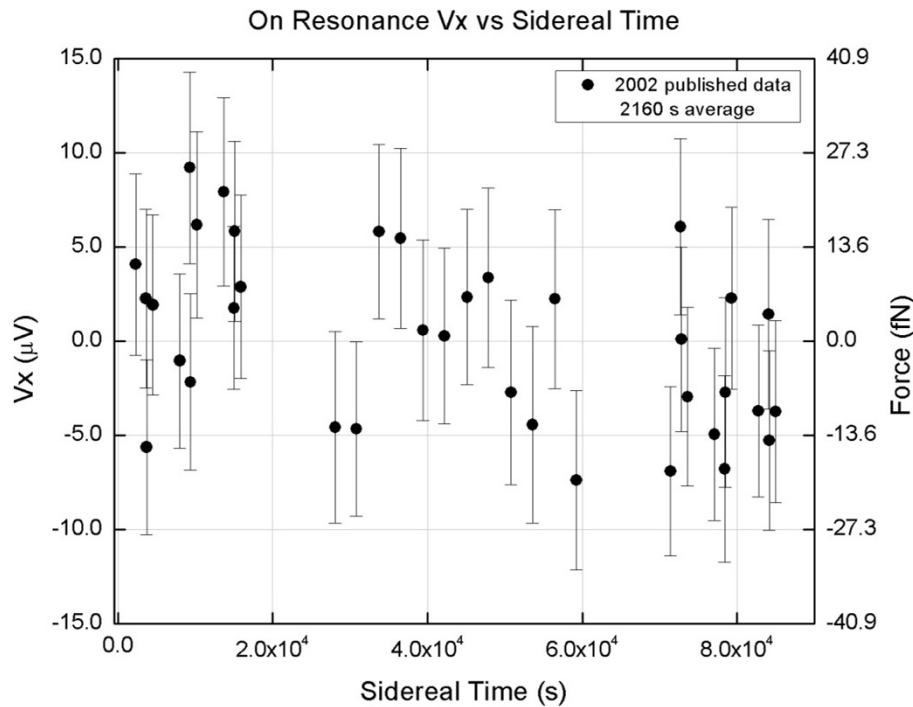
$$S_{LV} = f(u, \underbrace{s^{\mu\nu}, t^{\kappa\lambda\mu\nu}})$$

20 coefficients controlling L.V.

Estimated sensitivities:  $10^{-15} - 10^{-4}$

# 2002 Data as Function of Time

- 22 hrs of data accumulated over 5 days (August 2002)
- On-resonance (signal) data accumulated in 12 minute sets (off-resonance, diagnostic data in between)
- Plots:  
Average signal over 3 consecutive sets (best for viewing time distribution) with  $1\sigma$  error, vs mean time of the sets





# Calculation of the Fitting Function

- LV force function [1]:
$$dF^{\hat{j}} = Gdm_1dm_2 \left( -\frac{x^{\hat{j}}}{x^3} - \frac{3}{2} \frac{x^{\hat{j}}}{x^5} x^{\hat{j}} x^{\hat{k}} \bar{s}^{\hat{i}\hat{k}} + \frac{x^{\hat{k}} \bar{s}^{\hat{j}\hat{k}}}{x^3} \right)$$
$$\bar{s}^{\hat{j}\hat{k}} = \text{coefficients of Lorentz violation in the SME standard lab frame}$$
$$(x_L = \text{South}, y_L = \text{East}, z_L = \text{vertical})$$

# Calculation of the Fitting Function

- LV force function [1]:
$$dF^{\hat{j}} = Gdm_1dm_2 \left( -\frac{x^{\hat{j}}}{x^3} - \frac{3}{2} \frac{x^{\hat{j}}}{x^5} x^{\hat{j}} x^{\hat{k}} \bar{s}^{\hat{i}\hat{k}} + \frac{x^{\hat{k}} \bar{s}^{\hat{j}\hat{k}}}{x^3} \right)$$
$$\bar{s}^{\hat{j}\hat{k}} = \text{coefficients of Lorentz violation in the SME standard lab frame}$$
$$(x_L = \text{South}, y_L = \text{East}, z_L = \text{vertical})$$

**Force misaligned relative to  $\vec{r} = \vec{x}_1 - \vec{x}_2$ , but  $1/r^2$  behavior preserved**

# Calculation of the Fitting Function

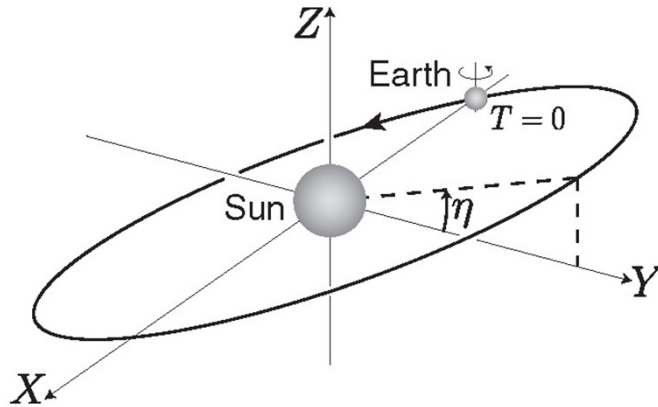
- LV force function [1]:
 
$$dF^{\hat{j}} = Gdm_1dm_2 \left( -\frac{x^{\hat{j}}}{x^3} - \frac{3}{2} \frac{x^{\hat{j}}}{x^5} x^{\hat{j}} x^{\hat{k}} \bar{s}^{\hat{i}\hat{k}} + \frac{x^{\hat{k}} \bar{s}^{\hat{j}\hat{k}}}{x^3} \right)$$

$$\bar{s}^{\hat{j}\hat{k}} = \text{coefficients of Lorentz violation in the SME standard lab frame}$$

( $x_L = \text{South}, y_L = \text{East}, z_L = \text{vertical}$ )

**Force misaligned relative to  $\vec{r} = \vec{x}_1 - \vec{x}_2$ , but  $1/r^2$  behavior preserved**

- Transform to sun-centered frame [2]:



$$R = \begin{pmatrix} \cos \chi \cos \omega_{\oplus} T & \cos \chi \sin \omega_{\oplus} T & -\sin \chi \\ -\sin \omega_{\oplus} T & \cos \omega_{\oplus} T & 0 \\ \sin \chi \cos \omega_{\oplus} T & \sin \chi \sin \omega_{\oplus} T & \cos \chi \end{pmatrix}$$

$\omega_{\oplus}$  = sidereal frequency

ignore boost ( $\eta$ );  $\chi = \text{colatitude} = 0.87$

[1] Q. G. Bailey and V. A. Kostelecký, PRD 74 045001 (2006).

[2] V. A. Kostelecký and M. Mewes, PRD 66 056005 (2002).

# Calculation of the Fitting Function

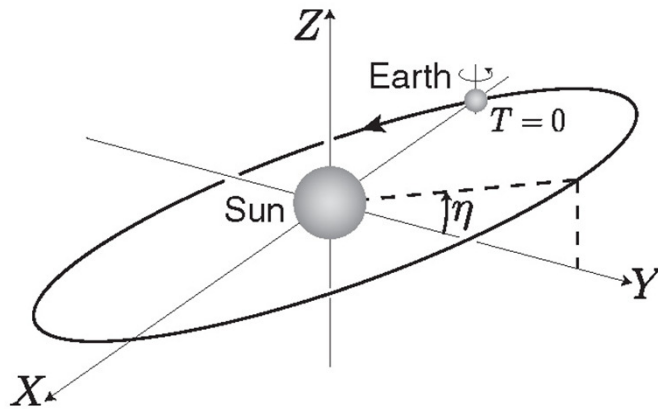
- LV force function [1]:
 
$$dF^{\hat{j}} = Gdm_1dm_2 \left( -\frac{x^{\hat{j}}}{x^3} - \frac{3}{2} \frac{x^{\hat{j}}}{x^5} x^{\hat{k}} x^{\hat{l}} \bar{s}^{\hat{i}\hat{k}} + \frac{x^{\hat{k}} \bar{s}^{\hat{j}\hat{k}}}{x^3} \right)$$

$$\bar{s}^{\hat{j}\hat{k}} = \text{coefficients of Lorentz violation in the SME standard lab frame}$$

( $x_L = \text{South}, y_L = \text{East}, z_L = \text{vertical}$ )

**Force misaligned relative to  $\vec{r} = \vec{x}_1 - \vec{x}_2$ , but  $1/r^2$  behavior preserved**

- Transform to sun-centered frame [2]:



$$R = \begin{pmatrix} \cos \chi \cos \omega_{\oplus} T & \cos \chi \sin \omega_{\oplus} T & -\sin \chi \\ -\sin \omega_{\oplus} T & \cos \omega_{\oplus} T & 0 \\ \sin \chi \cos \omega_{\oplus} T & \sin \chi \sin \omega_{\oplus} T & \cos \chi \end{pmatrix}$$

$\omega_{\oplus}$  = sidereal frequency

ignore boost ( $\eta$ );  $\chi = \text{colatitude} = 0.87$

- Detector has distributed mass:

$$F = \frac{1}{|z_{\max}|} \int_D d^3 \vec{x} \vec{z}^F(\vec{x}) \cdot dF^{\hat{j}}(\vec{x})$$

[1] Q. G. Bailey and V. A. Kostelecký, PRD 74 045001 (2006).

[2] V. A. Kostelecký and M. Mewes, PRD 66 056005 (2002).

# Calculation of the Fitting Function

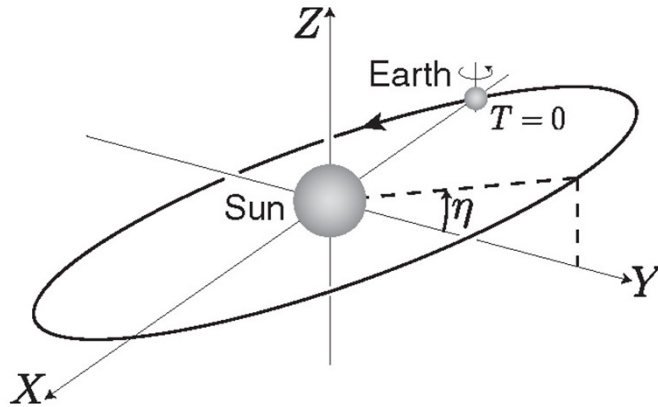
- LV force function [1]:
 
$$dF^{\hat{j}} = Gdm_1dm_2 \left( -\frac{x^{\hat{j}}}{x^3} - \frac{3}{2} \frac{x^{\hat{j}}}{x^5} x^{\hat{k}} x^{\hat{l}} \bar{s}^{\hat{i}\hat{k}} + \frac{x^{\hat{k}} \bar{s}^{\hat{j}\hat{k}}}{x^3} \right)$$

$$\bar{s}^{\hat{j}\hat{k}} = \text{coefficients of Lorentz violation in the SME standard lab frame}$$

( $x_L = \text{South}, y_L = \text{East}, z_L = \text{vertical}$ )

**Force misaligned relative to  $\vec{r} = \vec{x}_1 - \vec{x}_2$ , but  $1/r^2$  behavior preserved**

- Transform to sun-centered frame [2]:



$$R = \begin{pmatrix} \cos \chi \cos \omega_{\oplus} T & \cos \chi \sin \omega_{\oplus} T & -\sin \chi \\ -\sin \omega_{\oplus} T & \cos \omega_{\oplus} T & 0 \\ \sin \chi \cos \omega_{\oplus} T & \sin \chi \sin \omega_{\oplus} T & \cos \chi \end{pmatrix}$$

$\omega_{\oplus}$  = sidereal frequency

ignore boost ( $\eta$ );  $\chi$  = colatitude = 0.87

- Detector has distributed mass:

$$F = \frac{1}{|z_{\max}|} \int_D d^3 \vec{x} \vec{z}^F(\vec{x}) \cdot dF^{\hat{j}}(\vec{x})$$

mode shape from finite element model

[1] Q. G. Bailey and V. A. Kostelecký, PRD 74 045001 (2006).

[2] V. A. Kostelecký and M. Mewes, PRD 66 056005 (2002).

# Results

*D. Bennett, V. Skavysh, J. Long, Proc. 5<sup>th</sup> CPT conference*

- Force:

$$F = C_0 + S_\omega \sin(\omega_\oplus T) + C_\omega \cos(\omega_\oplus T) + S_{2\omega} \sin(2\omega_\oplus T) + C_{2\omega} \cos(2\omega_\oplus T)$$

$C_\omega, S_\omega$  functions of detector geometry,  $s^{JK}$

| Component     | Amplitude ( $\times 10^{-16}$ N)   |
|---------------|--|
| $C_0$         | $1.12\bar{s}^{XX} + 0.00\bar{s}^{XY} - 7.78\bar{s}^{XZ} - 3.48\bar{s}^{YY} + 0.00\bar{s}^{YZ} - 0.21\bar{s}^{ZZ}$  |
| $S_\omega$    | $0.07\bar{s}^{XX} + 0.76\bar{s}^{XY} + 0.02\bar{s}^{XZ} + 0.00\bar{s}^{YY} - 0.64\bar{s}^{YZ} - 0.07\bar{s}^{ZZ}$  |
| $C_\omega$    | $-0.49\bar{s}^{XX} + 0.11\bar{s}^{XY} - 0.17\bar{s}^{XZ} + 0.00\bar{s}^{YY} - 0.09\bar{s}^{YZ} + 0.49\bar{s}^{ZZ}$ |
| $S_{2\omega}$ | $0.06\bar{s}^{XX} - 0.08\bar{s}^{XY} + 0.15\bar{s}^{XZ} - 0.16\bar{s}^{YY} - 0.10\bar{s}^{YZ} - 0.09\bar{s}^{ZZ}$  |
| $C_{2\omega}$ | $0.03\bar{s}^{XX} + 0.20\bar{s}^{XY} + 0.06\bar{s}^{XZ} - 0.06\bar{s}^{YY} + 0.24\bar{s}^{YZ} + 0.04\bar{s}^{ZZ}$  |

- Fit:

| Coefficient    | Mean and error ( $2\sigma$ )   |
|----------------|--------------------------------|
| $\bar{s}^{XX}$ | $(-0.04 \pm 4.90) \times 10^4$ |
| $\bar{s}^{XY}$ | $(-0.07 \pm 6.12) \times 10^4$ |
| $\bar{s}^{XZ}$ | $(-0.01 \pm 2.56) \times 10^3$ |
| $\bar{s}^{YZ}$ | $(-0.06 \pm 5.83) \times 10^4$ |
| $\bar{s}^{ZZ}$ | $(0.08 \pm 6.68) \times 10^4$  |

Compare: Chung, Chu, et al., PRD 80 016002:  $S^{JK} < 1 \times 10^{-8}$   
 (atom interferometer sensitive to  $\Delta g/g \sim 1 \times 10^{-9}$ )