



Shining light on dark energy and modifications of gravity

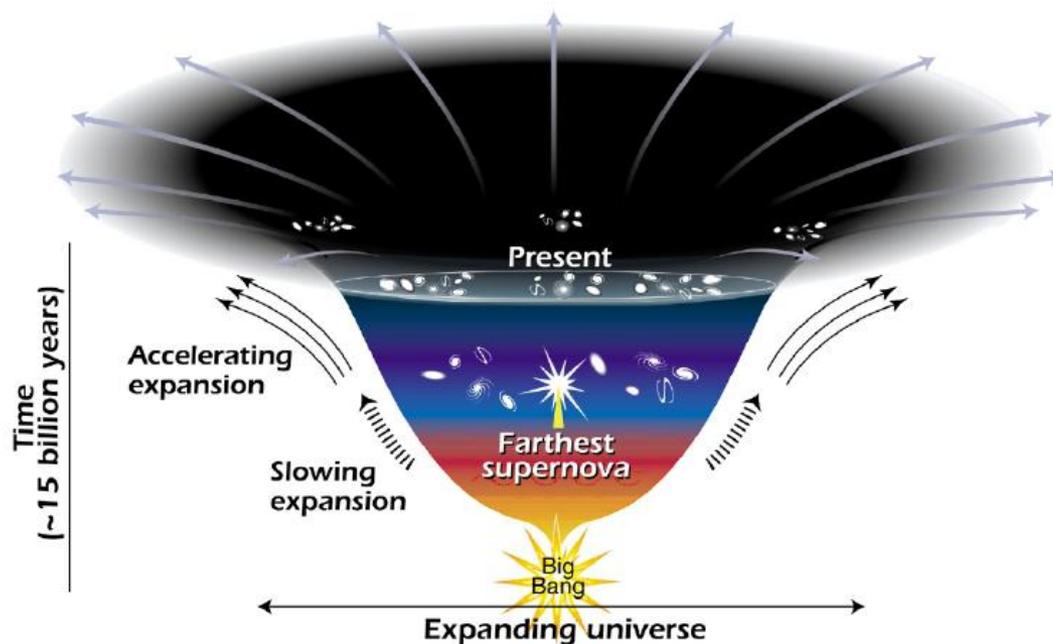
Clare Burrage
University of Nottingham

With Philippe Brax and Anne Davis

arXiv:1206.1809

Dark energy and modified gravity

- Almost all proposals directly or indirectly introduce a new light scalar field



- The cause of the acceleration of the expansion of the universe must be coherent across the universe
 - This corresponds to very light masses

What is the most general way of interacting with matter?

- Matter fields couple to an effective metric which depends on the scalar field

$$\mathcal{L} \supset \tilde{g}_{\mu\nu} T^{\mu\nu}$$

- The most general effective metric that can be made from the metric and a scalar field that respects causality and the weak equivalence principle is

$$\tilde{g}_{\mu\nu} = A(\phi, X)g_{\mu\nu} + B(\phi, X)\partial_\mu\phi\partial_\nu\phi$$



Conformal



Disformal

$$X = -(\partial\phi)^2/2.$$

- Disformal terms don't give rise to fifth forces

(Bekenstein 1993)

Motivation for disformal terms

- Massive gravity (de Rham, Gabadadze 2010)
 - A massive graviton can be decomposed into
 - One helicity-two mode
 - Two helicity-one modes
 - One helicity-zero mode
- In ghost free formulations of massive gravity matter fields couple to a metric

$$h_{\mu\nu} = \hat{h}_{\mu\nu} + \pi\eta_{\mu\nu} + \frac{(6c_3 - 1)}{\Lambda_3^3} \partial_\mu \pi \partial_\nu \pi$$

- Where the scale is $\Lambda_3^3 = M_P m^2$,
- Also disformal terms from: quintessence, inflation, varying speed of light...

Energy scales

- A dark energy motivation leads us to expect

$$m \sim H_0 \sim 10^{-33} \text{ eV}$$

- A modified gravity motivation leads us to expect

$$m \sim H_0 \sim 10^{-33} \text{ eV}$$

$$M \sim \Lambda \sim M_P \sim 10^{19} \text{ GeV}$$

- Massive gravity gives

$$m \sim H_0 \sim 10^{-33} \text{ eV}$$

$$\Lambda \sim M_P \sim 10^{19} \text{ GeV}$$

$$M \sim \sqrt{M_P m_{\text{grav}}} \sim 10^{-1} \text{ eV}$$

The Vainshtein Mechanism

- The scalar field from massive gravity has higher order derivative self interactions
- A massive object gives rise to a non-trivial scalar field configuration
- The self interaction terms mean that the coupling constants for perturbations around the background vary

$$\tilde{\Lambda} = Z\Lambda$$

$$\tilde{M} = Z^{1/2}M$$

$$Z \gg 1$$

A photograph of a large, arched, multi-paned window. The window is filled with a bright, warm yellow light, creating a strong contrast with the dark background. A beam of light extends from the window towards the right side of the frame. The window has a grid pattern of panes, with a smaller arched section at the top.

Interactions with photons

Scalar Fields Couple to Gauge Bosons

- Conformally coupled scalar fields

$$\tilde{g}_{\mu\nu} = A(\phi)g_{\mu\nu}$$

- There is no coupling to the kinetic terms of gauge bosons at tree level

$$\mathcal{L} \supset \sqrt{-g} g^{\mu\nu} g^{\alpha\beta} F_{\mu\alpha} F_{\nu\beta}$$

- The Standard Model is not conformally invariant
 - The conformal anomaly means that a coupling to photons will always be generated
 - The scale of coupling is undetermined

Interactions with photons

- The general Lagrangian for scalar fields and photons is

$$\mathcal{L}_{\phi,\gamma} = -\frac{1}{2}(\partial\phi)^2 - V(\phi) - \frac{1}{4}F^2 - \frac{\phi}{\Lambda}F^2 - \frac{1}{M^4}\partial_\mu\phi\partial_\nu\phi \left[\frac{1}{4}g^{\mu\nu}F^2 - F^\mu{}_\alpha F^{\nu\alpha} \right]$$

- Affects the propagation of photons through magnetic fields, causing
 - Changes in polarization
 - Photon number non-conservation
- Then the relevance of the disformal term is controlled by

$$b = \frac{B}{M^2}$$

Propogation through a magnetic field

- This system can be diagonalised and solved

$$\begin{pmatrix} \phi(x) \\ A_y(x) \end{pmatrix} = P \begin{pmatrix} e^{-i\omega(1+\lambda_+)x} & 0 \\ 0 & e^{-i\omega(1+\lambda_-)x} \end{pmatrix} P^{-1} \begin{pmatrix} \phi(0) \\ A_y(0) \end{pmatrix}$$

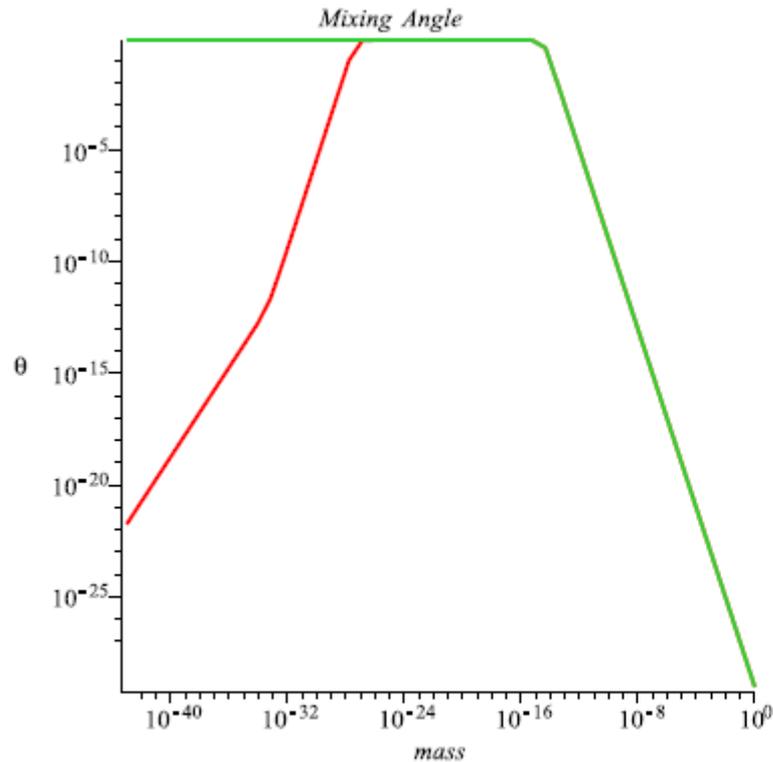
- The mixing matrix is

$$P = \begin{pmatrix} \sin \vartheta & -\cos \vartheta \\ \cos \vartheta & \sin \vartheta \end{pmatrix} \quad \tan 2\vartheta = \frac{4B}{\Lambda\omega} \sqrt{\frac{1+b^2}{1-a^2}} \left(\frac{m^2}{2\omega^2} - b^2 \right)^{-1}$$

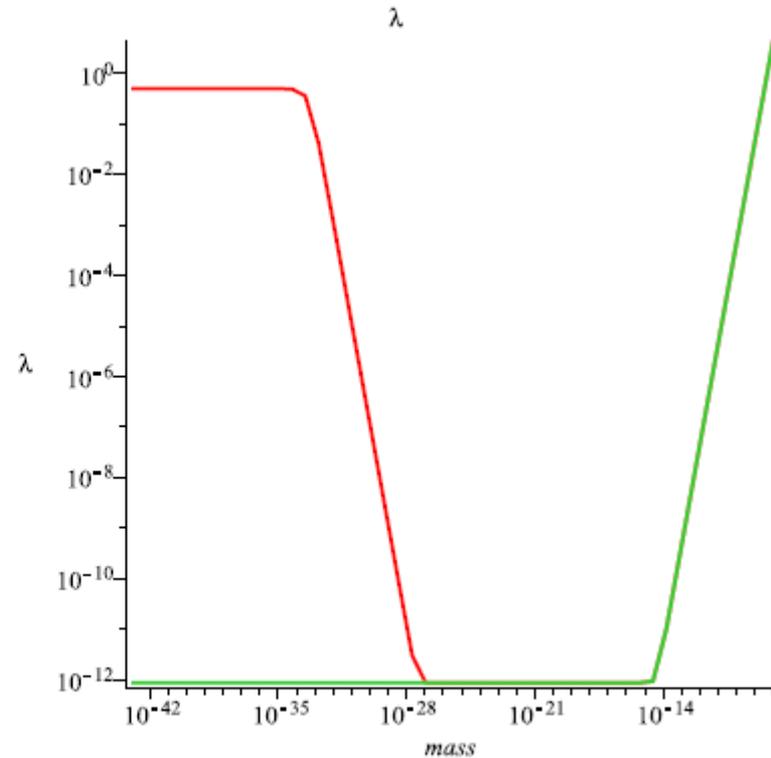
- The eigenvalues are

$$\lambda_{\pm} = -\lambda(\cos 2\vartheta \mp 1) \quad \lambda = \frac{1}{2(1+b^2)} \left| \frac{m^2}{2\omega^2} - b^2 \right| (1 + \tan^2 2\vartheta)^{1/2}$$

Propogation through a magnetic field



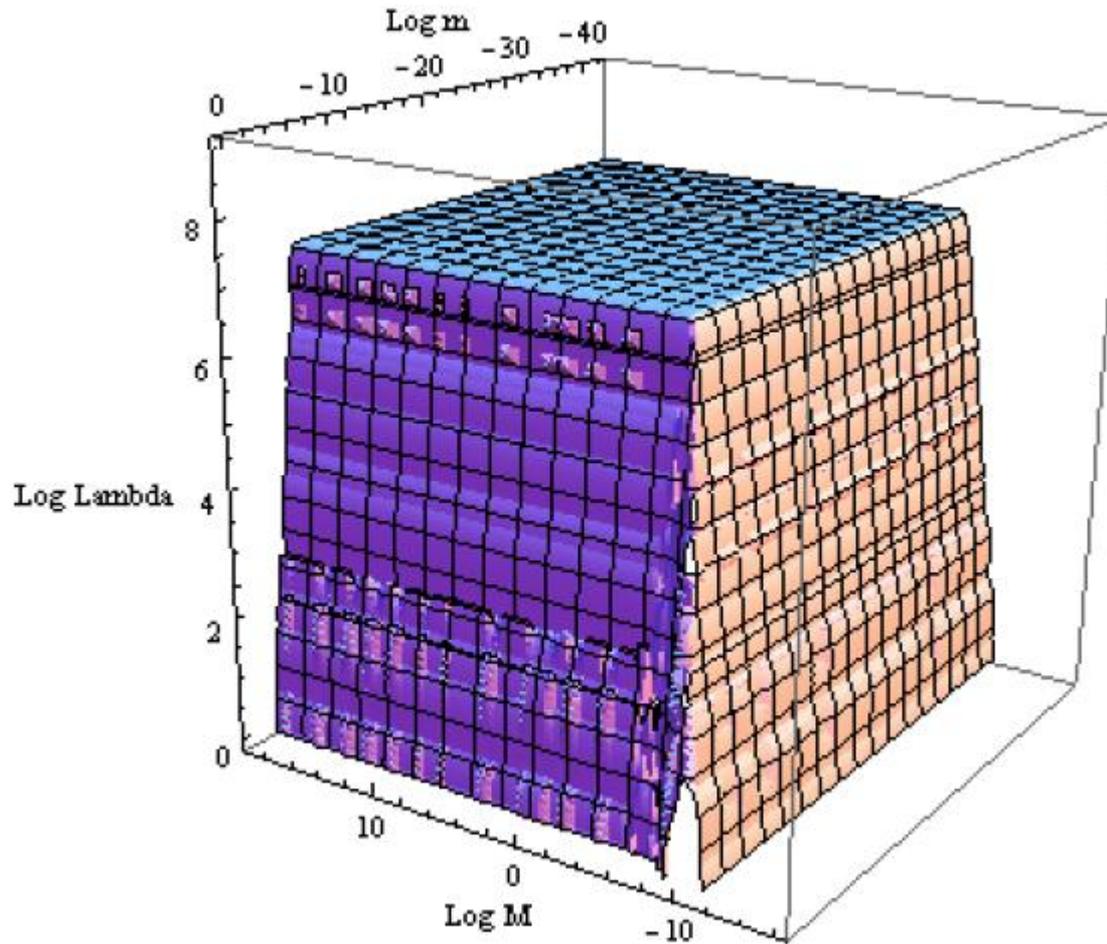
(a) Mixing Angle, ϑ



(b) Oscillation Wavelength, λ

$$\Lambda = 10^6 \text{ GeV} \quad M^2 = m M_P \quad B = 5 \text{ Tesla} \quad \omega = 2.33 \text{ eV}$$

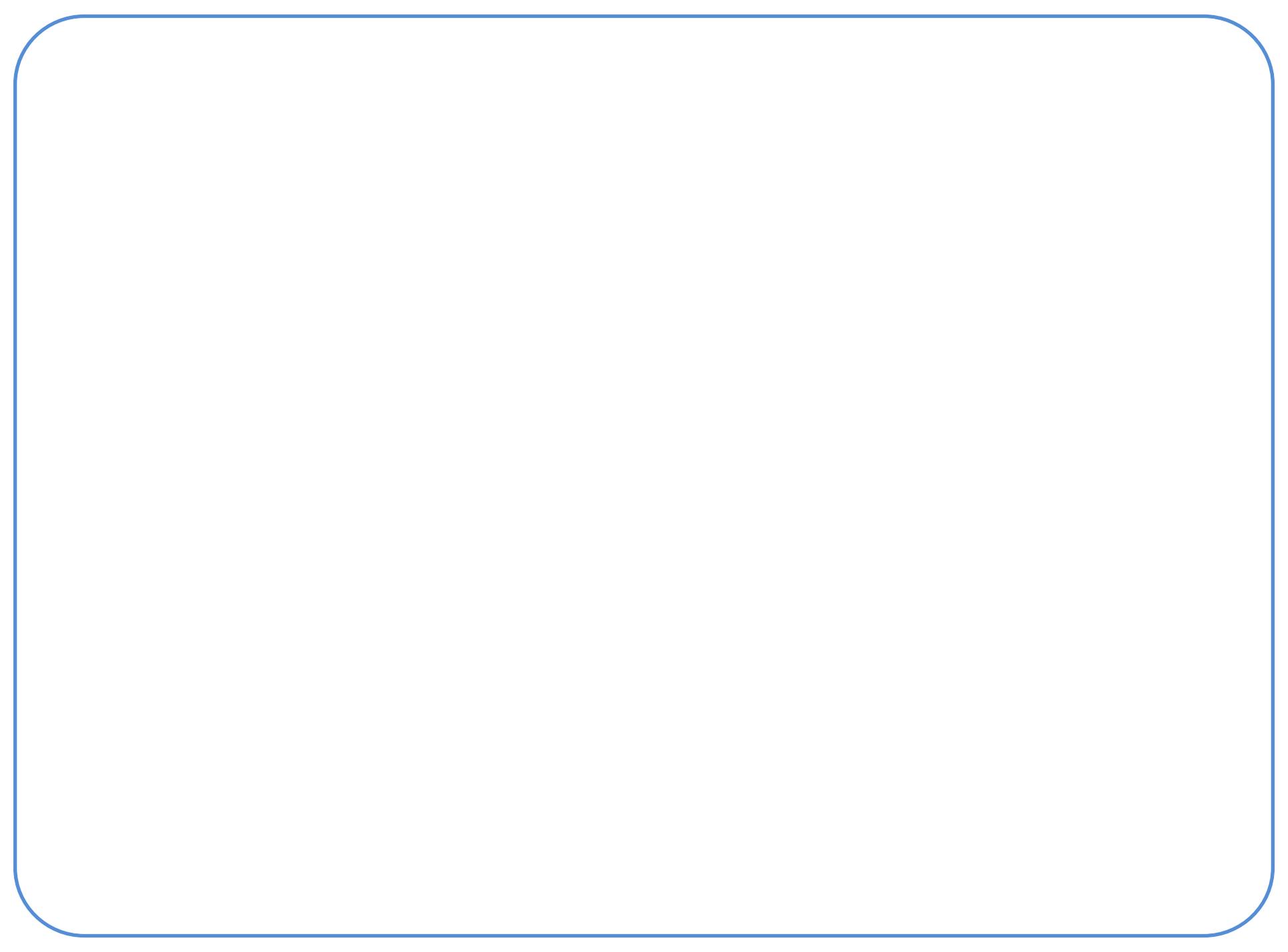
Constraints of the ALPS experiment



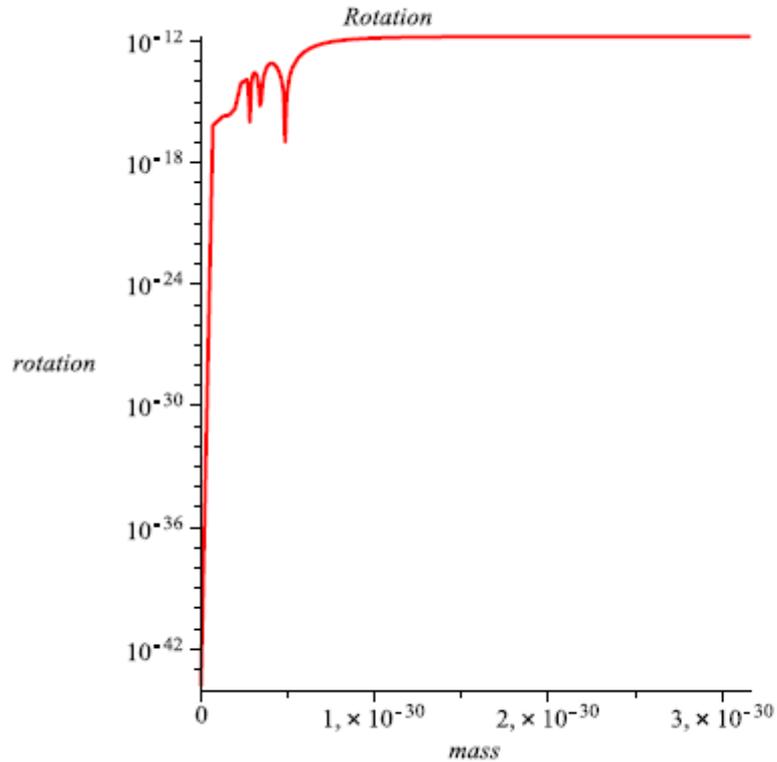
ALPS results from Ehret et al. 2010

Conclusions

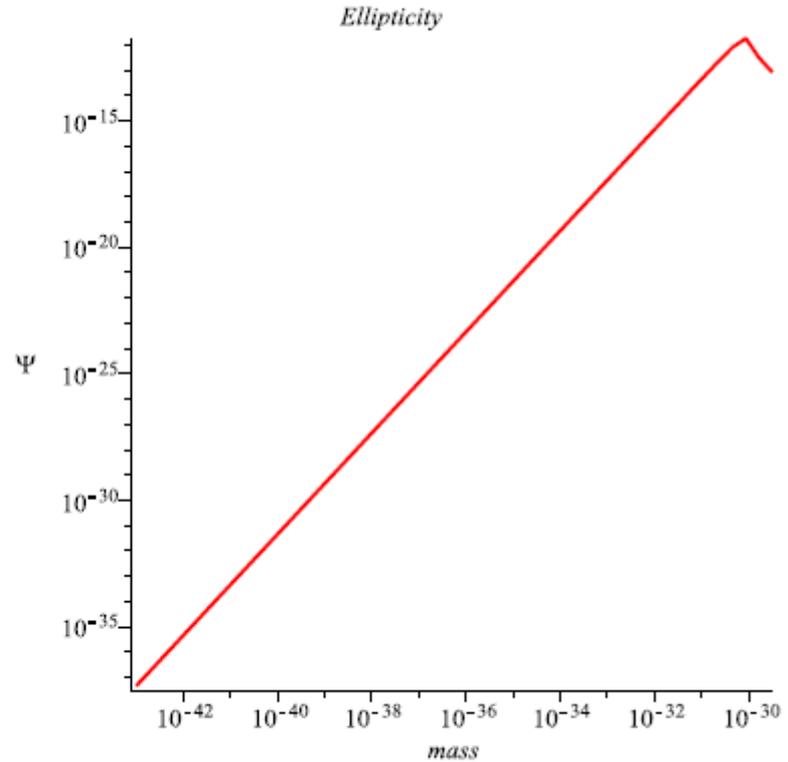
- Dark energy and modifications of gravity introduce new scalar degrees of freedom
- The most general coupling includes conformal and disformal terms
- The disformal terms can couple at a low energy scale
 - eg massive gravity
- High-precision low-energy photon experiments can be used to search for the scalar components of massive gravitons



Scalars at BMV



(a) Rotation

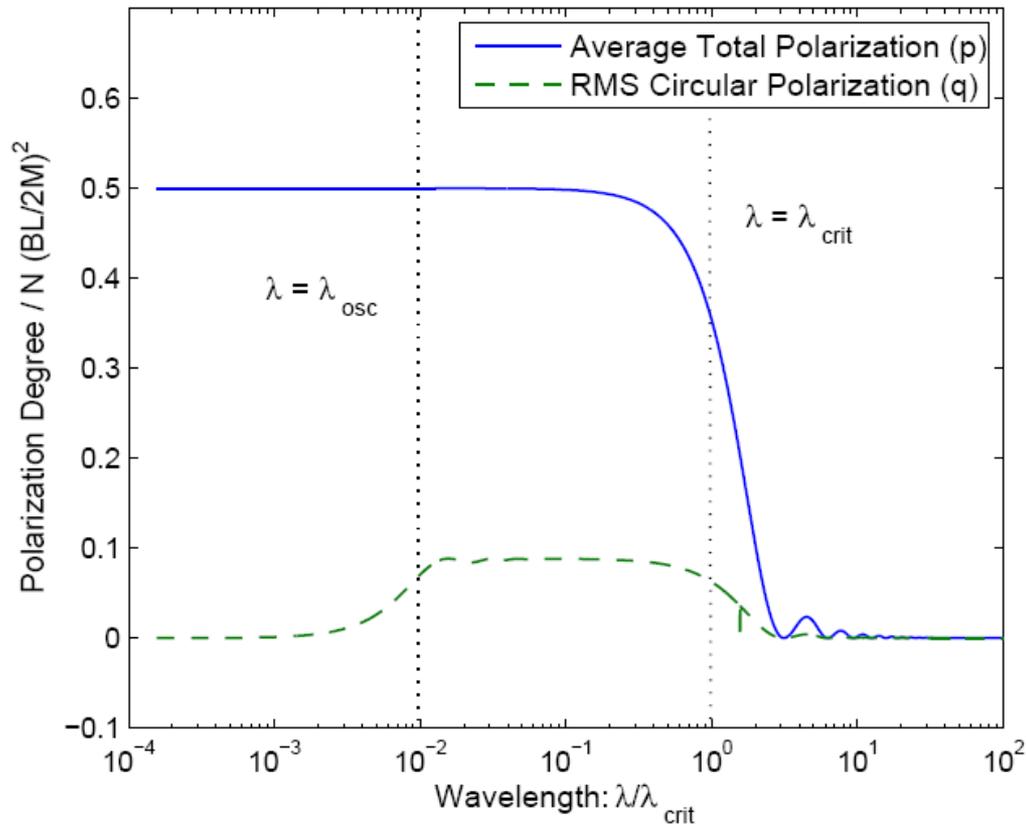


(b) Ellipticity

$$\Lambda = 10^6 \text{ GeV} \quad \phi_0 = 10^{-2} \Lambda \quad M^2 = m M_P \quad B = 9 \text{ Tesla} \quad \omega = 1.17 \text{ eV}$$

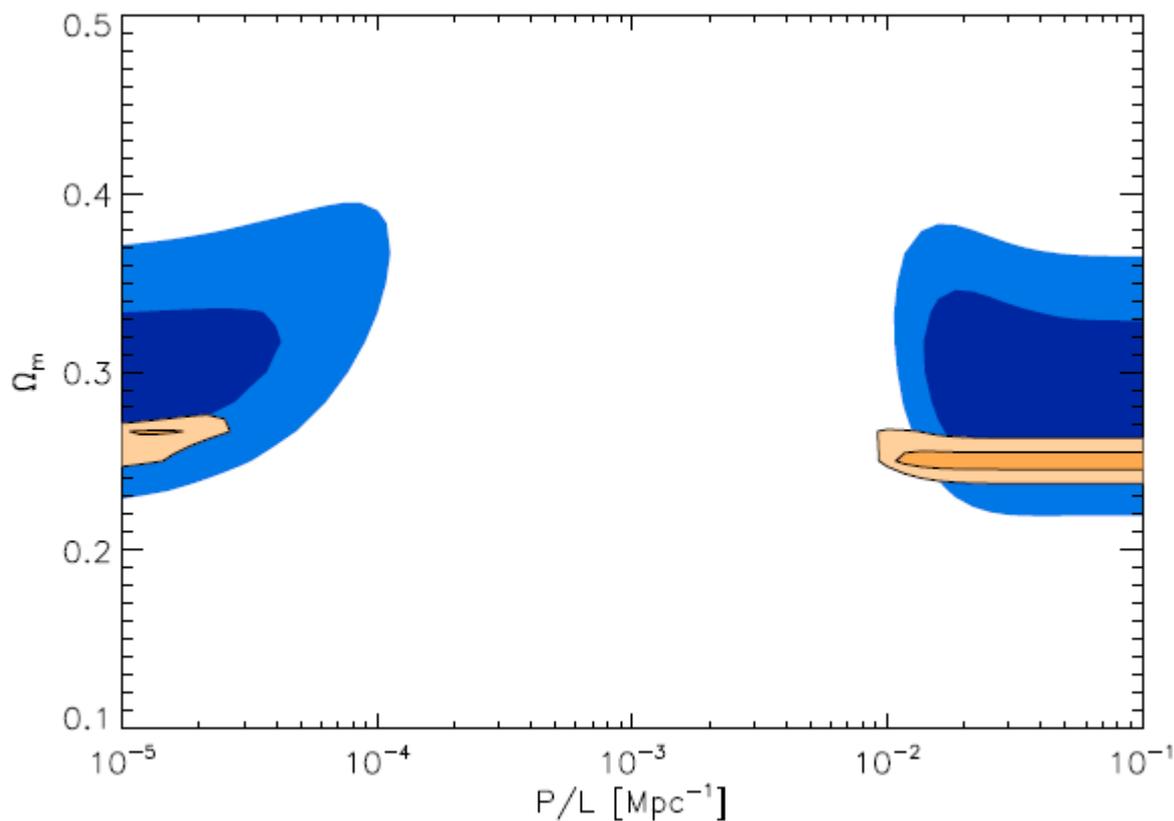
Starlight polarisation

- For a conformal scalar:

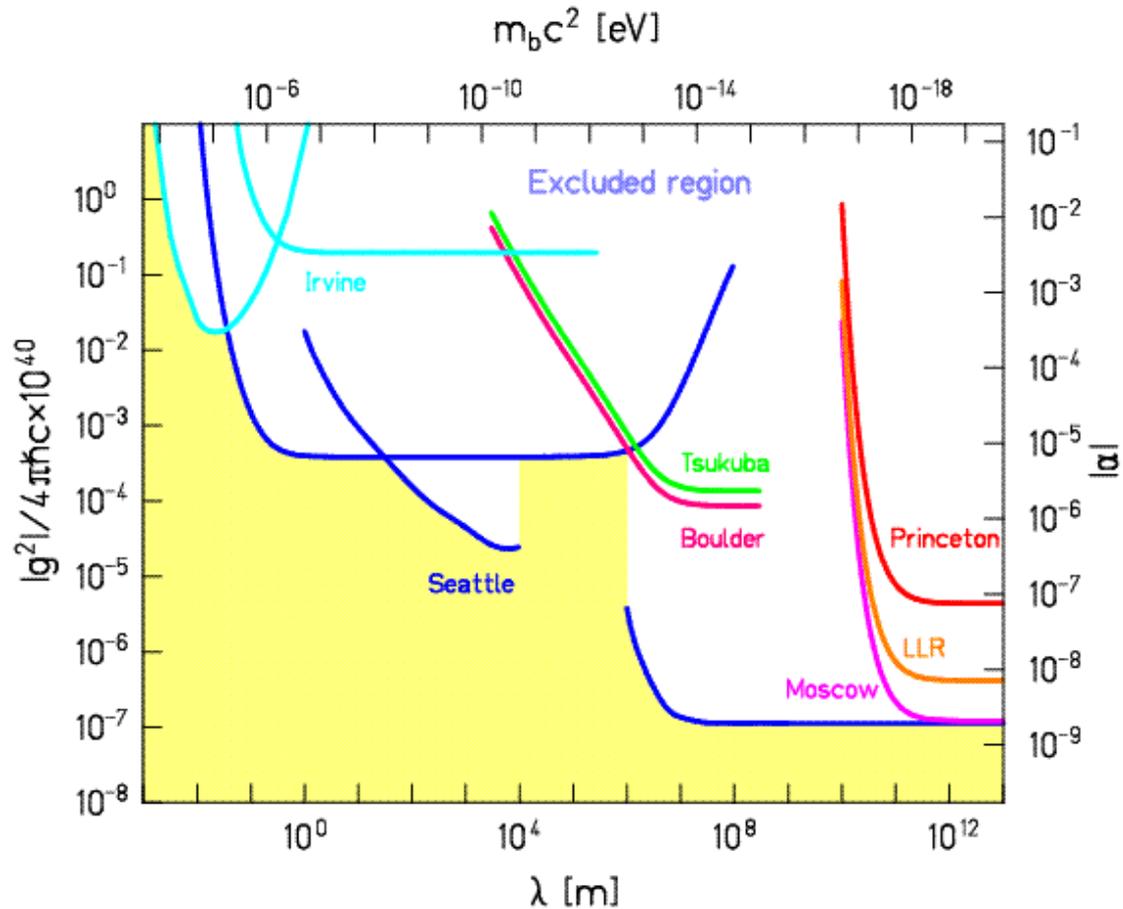


Cosmic Opacity

- Violations of photons conservation will look like extra opacity in the universe
- For a conformal scalar



95% confidence limits on
Equivalence Principle violating Yukawa interactions
coupled to baryon number



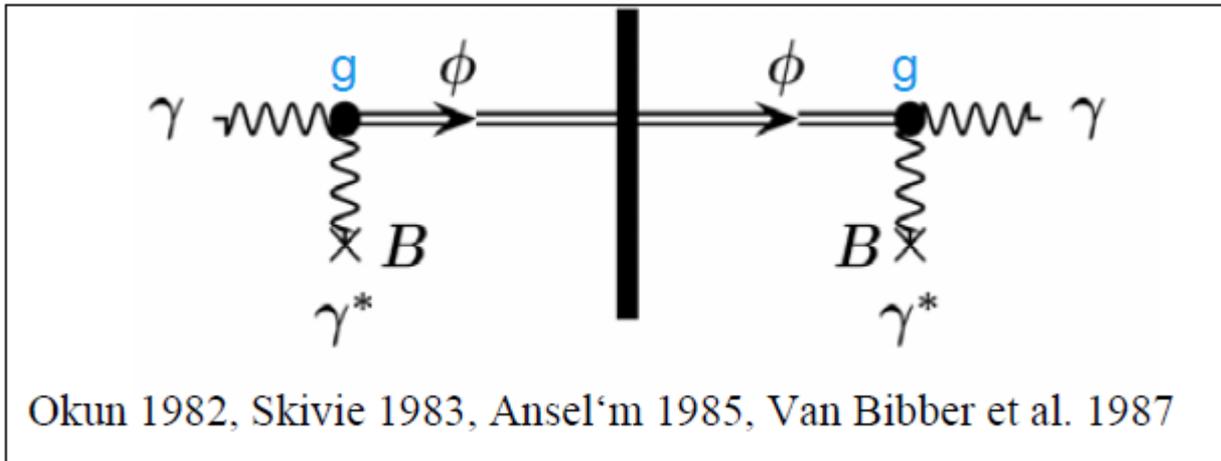
Results of the Eöt-Wash experiment at the University of Washington

Light shining through walls

- The scalar coupling

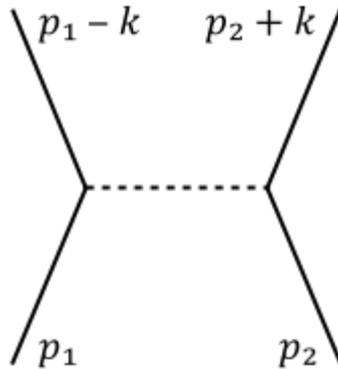
$$\mathcal{L} \supset \frac{\phi}{M_\gamma} F_{\mu\nu} F^{\mu\nu}$$

- Leads to light shining through walls



Fifth forces

- Light scalar fields mediate new fifth forces



- Forces with conformal couplings seem to be in direct conflict with experimental searches
- To include them we have to either
 - Suppress their coupling to matter with an energy scale above the Planck scale
 - Make the theory non-linear eg chameleons

Fifth forces

- A purely disformal coupling

$$\tilde{g}_{\mu\nu} = g_{\mu\nu} + \frac{2\partial_\mu\phi\partial_\nu\phi}{M^4}$$
$$\mathcal{L}_\phi = -\frac{1}{2}(\partial\phi)^2 - \frac{1}{2}m^2\phi^2 + \frac{1}{M^4}\partial_\mu\phi\partial_\nu\phi T^{\mu\nu}$$

- Gives rise to equations of motion

$$\square\phi - m^2\phi - \frac{2}{M^4}\nabla_\mu(\partial_\nu\phi T^{\mu\nu}) = 0$$

- A static non-relativistic object does not source a scalar field profile, so there are no fifth forces

Fifth forces

- With conformal and disformal terms

$$\tilde{g}_{\mu\nu} = \left(1 + \frac{\phi}{\Lambda}\right) g_{\mu\nu} + \frac{2}{M^4} \partial_\mu \phi \partial_\nu \phi$$

- The equation of motion is

$$\square\phi - m^2\phi + \frac{T}{\Lambda} - \frac{2}{M^4} \nabla_\mu (\partial_\nu \phi T^{\mu\nu}) = 0$$

- The conformal coupling sources a scalar field profile

$$\phi \approx \frac{M_c}{\Lambda r}$$

Fifth forces

- To compute the force we need the geodesic equation

$$\tilde{u}^\nu \tilde{\nabla}_\nu \tilde{u}^\mu = 0 \qquad \tilde{g}_{\mu\nu} \tilde{u}^\mu \tilde{u}^\nu = -1$$

- We can rewrite this with quantities defined wrt the other metric

$$g_{\mu\nu} u^\mu u^\nu = -1 \qquad a^\mu = u^\nu \nabla_\nu u^\mu$$

$$a^\mu = F^\mu(\phi, \partial\phi, \partial\partial\phi)$$

- For a static, spherically symmetric, non-relativistic source

$$F_r = -\frac{\phi'}{2\Lambda} \left(1 + \frac{\phi'^2}{M^4 + \phi'^2} \right)$$

Motivation for disformal terms

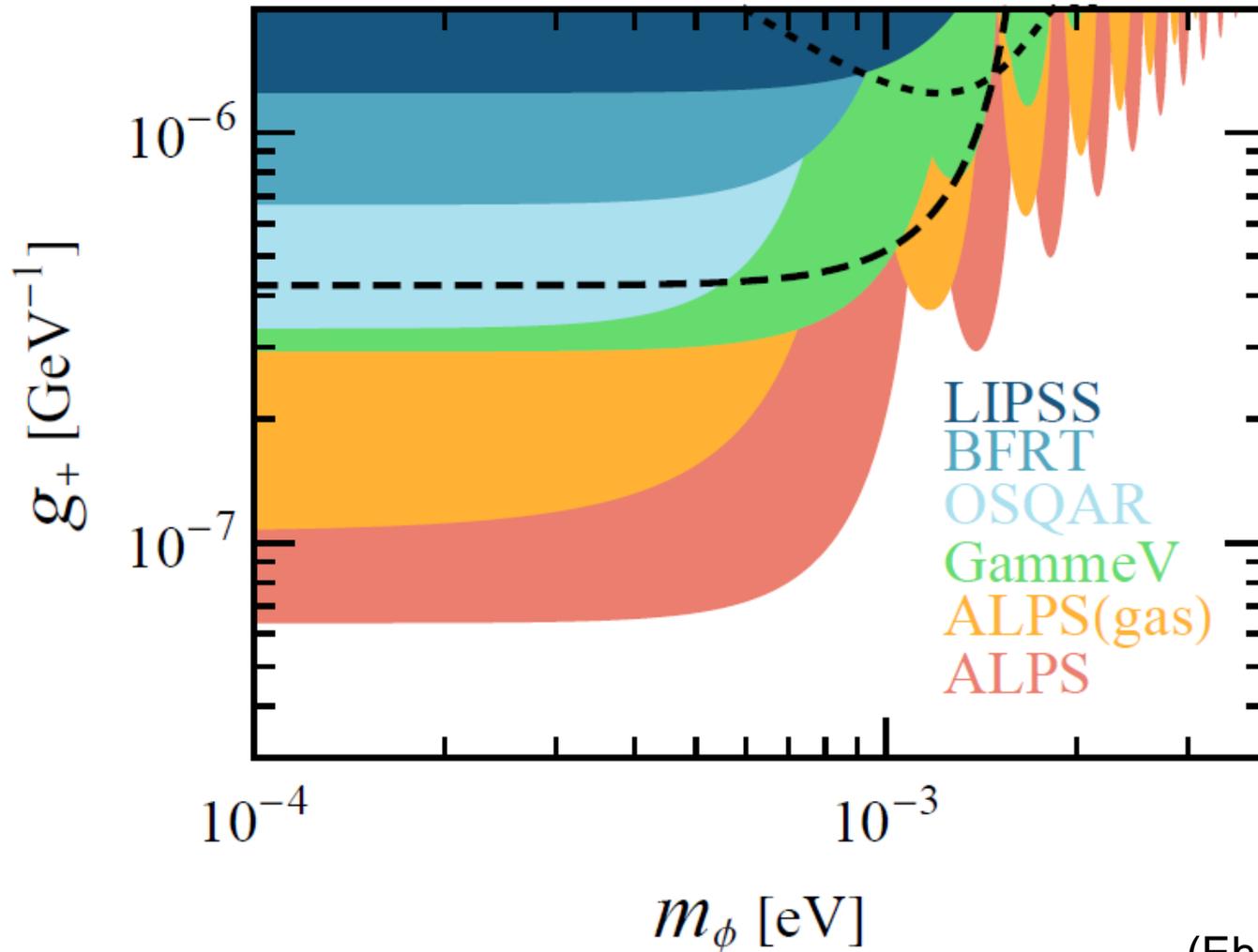
- Disformal inflation
 - Allows for a graceful exit but requires a curvaton

(Kaloper 2004)
- Disformal quintessence
 - Helps alleviate the coincidence problem
 - No fifth force problems

(Koivisto 2008. Kovisto, Mota, Zumalacárregui 2012)
- Varying speed of light
 - Alternative explanations for the apparent acceleration of the expansion of the universe
 - Now some tension with observations

(Clayton, Moffat 2000. Drummond 1999. Magueijo 2003)

Constraints on the conformal coupling



(Ehret et al. 2010)

Disformal mixing with photons

$$\mathcal{L}_{\phi,\gamma} = -\frac{1}{2}(\partial\phi)^2 - V(\phi) - \frac{1}{4}F^2 - \frac{\phi}{\Lambda}F^2 - \frac{1}{M^4}\partial_\mu\phi\partial_\nu\phi \left[\frac{1}{4}g^{\mu\nu}F^2 - F^\mu{}_\alpha F^{\nu\alpha} \right]$$

- The background solutions

$$A_0 = 0 ,$$

$$A_i = \frac{1}{2}\epsilon_{ijk}B_jx_k$$

$$V'(\phi_0) = -\frac{2}{\Lambda}B^2$$

- We want to study perturbations around this background

$$\phi \rightarrow \phi_0 + \phi ,$$

$$A_\mu \rightarrow \frac{1}{2}\delta_{\mu i}\epsilon_{ijk}B_jx_k + A_\mu .$$

Propogation through a magnetic field

- Equations of motion

$$\square\phi \left(1 + \frac{B^2}{M^4}\right) - \frac{2}{M^4}(\nabla\phi B^2 - \partial_i\partial_j\phi B^i B^j) = m^2\phi - \frac{2}{\Lambda}\epsilon_{ijk}B_j(\partial_k A_i - \partial_i A_k)$$

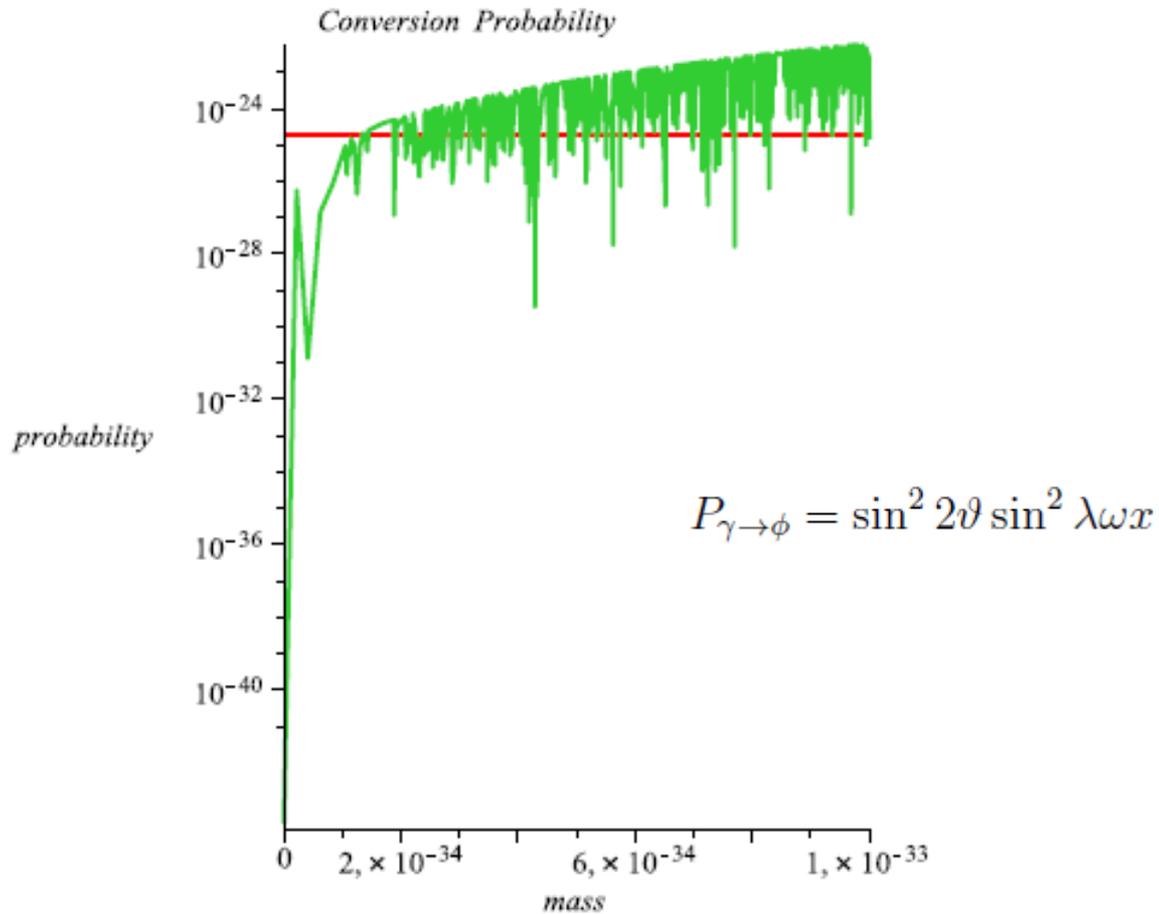
$$\left(1 + \frac{4\phi_0}{\Lambda}\right)\square A_\mu + \frac{4}{\Lambda}\delta_{\mu i}B_j\epsilon_{ijk}\partial_k\phi = 0$$

- Only one component of the photon mixes with the scalar

$$\begin{pmatrix} (\omega^2 - k^2) + \frac{2k^2b^2 - m^2}{1+b^2} & \frac{4Bk}{\Lambda\sqrt{1-a^2}\sqrt{1+b^2}} \\ \frac{4Bk}{\Lambda\sqrt{1+b^2}\sqrt{1-a^2}} & (\omega^2 - k^2) \end{pmatrix} \begin{pmatrix} \phi \\ A_y \end{pmatrix} = 0$$

$$a = 2\sqrt{\frac{-\phi_0}{\Lambda}} \quad b = \frac{B}{M^2}$$

Probability of mixing



$$\Lambda = 10^6 \text{ GeV} \quad M^2 = m M_P \quad B = 5 \text{ Tesla} \quad \omega = 2.33 \text{ eV}$$